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A NEW LEVINSON'S THEOREM FOR POTENTIALS WITH CRITICAL DECAY

XIAOYAO JIA, FRANÇOIS NICOLEAU, AND XUE PING WANG

ABSTRACT. We study the low-energy asymptotics of the spectral shift function for Schrödinger operators with potentials decaying like $O(\frac{1}{|x|^2})$. We prove a generalized Levinson's for this class of potentials in presence of zero eigenvalue and zero resonance.

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1. INTRODUCTION

Threshold spectral analysis of Schrödinger operators plays an important role in many problems arising from non-relativistic quantum mechanics, such as low-energy scattering, propagation of cold neutrons in a crystal or the Efimov effect in many-particle systems (cf. [5, 23, 31]). Presence of zero resonance of Schrödinger operators is responsible for several striking physical phenomena. The Levinson's theorem which relates the phase shift to the number of eigenvalues and the zero resonance is one of the oldest topics in this domain (cf. [18, 20, 23]). Threshold spectral analysis is usually carried out for potentials decaying faster than $\frac{1}{|x|^2}$ (cf. [12, 21, 30]). Potentials with the critical decay $\frac{1}{|x|^2}$ appear in many interesting situations such as spectral analysis on manifolds with conical end or ion-atom scattering for N -body Schrödinger operators. The threshold spectral properties for potentials with critical decay are quite different from those of more quickly decaying potentials and the contribution of zero resonance and zero eigenvalues to the asymptotics of the resolvent at low-energy and to the long-time expansion of wave functions are studied in [6, 32, 33]. In this work, we shall study the low-energy asymptotics of some spectral shift function and prove a generalized Levinson's theorem for potentials with critical decay in presence of zero eigenvalue and zero resonance.

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Recall that for a pair of selfadjoint operators (P, P_0) on some Hilbert space \mathcal{H} , if

$$(P + i)^{-k} - (P_0 + i)^{-k} \text{ is of trace class for some } k \in \mathbb{N}^*, \quad (1.1)$$

the spectral shift function $\xi(\lambda)$ is defined as distribution on \mathbb{R} by

$$\text{Tr} (f(H) - f(H_0)) = - \int_{\mathbb{R}} f'(\lambda) \xi(\lambda) d\lambda, \quad \forall f \in \mathcal{S}(\mathbb{R}). \quad (1.2)$$

In fact, this relation only defines $\xi(\cdot)$ up to an additive constant. It can be fixed by, for example, requiring $\xi(\lambda) = 0$ for λ sufficiently negative if P and P_0 are both bounded from below.

In this paper, we are interested in the threshold spectral analysis of the Schrödinger operator $P = -\Delta + v(x)$ on $L^2(\mathbb{R}^n)$ with $n \geq 2$ and $v(x)$ a real function satisfying

$$v(x) = \frac{q(\theta)}{r^2} + O(\langle x \rangle^{-\rho_0}), \quad |x| \gg 1, \quad (1.3)$$

for some $q \in C(\mathbb{S}^{n-1})$ and $\rho_0 > 2$. Here $x = r\theta$ with $r = |x|$ and $\theta = \frac{x}{|x|}$. \mathbb{S}^{n-1} is the unit sphere. Let $-\Delta_{\mathbb{S}^{n-1}}$ denote the Laplace-Beltrami operator on \mathbb{S}^{n-1} . We assume throughout this paper that the smallest eigenvalue λ_1 of $-\Delta_{\mathbb{S}^{n-1}} + q(\theta)$ verifies

$$\lambda_1 > -\frac{1}{4}(n-2)^2. \quad (1.4)$$

This assumption ensures that the form associated to $\tilde{P}_0 = -\Delta + \frac{q(\theta)}{r^2}$ on $C_0^\infty(\mathbb{R}^n \setminus \{0\})$ is positive. We still denote by \tilde{P}_0 its Friedrich's realization as selfadjoint operator on $L^2(\mathbb{R}^n)$. Notice that the condition (1.1) is not satisfied for the pair $(P, -\Delta)$ by lack of decay on v , nor for the pair (P, \tilde{P}_0) because of the critical singularity at 0. This leads us to introduce a model operator P_0 in the following way. Let $0 \leq \chi_j \leq 1$ ($j = 1, 2$) be smooth functions on \mathbb{R}^n such that $\text{supp } \chi_1 \subset B(0, R_1)$, $\chi_1(x) = 1$ when $|x| < R_0$ and

$$\chi_1(x)^2 + \chi_2(x)^2 = 1.$$

Set

$$P_0 = \chi_1(-\Delta)\chi_1 + \chi_2\tilde{P}_0\chi_2,$$

on $L^2(\mathbb{R}^n)$. When $q(\theta) = 0$, we can take $P_0 = -\Delta$ and the main results of this work still hold. Under the assumption (1.4), one has $P_0 \geq 0$. The operator P will mainly be considered as a perturbation of P_0 . Under the condition that $v(x)$ is bounded and satisfies (1.3) for some $\rho_0 > n$, (1.1) is satisfied for $k > \frac{n}{2}$ and the spectral shift function $\xi(\lambda)$ for the pair (P, P_0) is well defined by (1.1).

High energy asymptotics of the spectral shift function of Schrödinger operators has been studied by many authors (see for example [1],[24],[25],[26],[35]). In particular, D. Robert proved in [25] that for the pair (P, P_0) introduced as above, if the potential v is smooth and satisfies

$$|\partial_x^\alpha (v(x) - \frac{q(\theta)}{r^2})| \leq C_\alpha \langle x \rangle^{-\rho_0 - |\alpha|}, \quad \text{for } |x| \text{ large.} \quad (1.5)$$

for some $\rho_0 > n$, then one has

- (i). $\xi(\lambda)$ is C^∞ in $(0, \infty)$;
- (ii). $\frac{d^k}{d\lambda^k} \xi(\lambda)$ has a complete asymptotic expansion for $\lambda \rightarrow \infty$,

$$\frac{d^k}{d\lambda^k} \xi(\lambda) \sim \lambda^{n/2-k-1} \sum_{j \geq 0} \alpha_j^{(k)} \lambda^{-j}.$$

We are mainly interested in the low-energy asymptotics of the spectral shift function $\xi(\cdot)$ for (P, P_0) . Although $\xi(\cdot)$ depends on the choice of cut-offs χ_1 and χ_2 , we shall see that physically interesting quantities in low-energy limit are independent of such choices.

The low-energy resolvent asymptotics for P is studied in [32, 33] for v of the form $v(x) = \frac{q(\theta)}{r^2} + w(x)$ with $w(x)$ bounded and satisfying $w(x) = O(\langle x \rangle^{-\rho_0})$ for some $\rho_0 > 2$. In this work, in order to define the spectral shift function, we assume $v(x)$ to be bounded and in the last section on Levinson's theorem we even need to assume it smooth. Therefore, we need to slightly modify the proofs of [33] to suit with the present situation.

The eigenvalues of the operator $-\Delta_{\mathbb{S}^{n-1}} + q(\theta)$ play an important role in the threshold spectral analysis of P (cf. [6, 32, 33]). Let

$$\sigma_\infty = \left\{ \nu; \nu = \sqrt{\lambda + \frac{(n-2)^2}{4}}, \lambda \in \sigma(-\Delta_{\mathbb{S}^{n-1}} + q(\theta)) \right\}. \quad (1.6)$$

Denote

$$\sigma_k = \sigma_\infty \cap]0, k], \quad k \in \mathbb{N}^*. \quad (1.7)$$

σ_1 is closely related to the properties of zero resonance. We say that 0 is a resonance of P if the equation $Pu = 0$ admits a solution u such that $\langle x \rangle^{-1}u \in L^2$, but $u \notin L^2$. u is then called a resonant state of P . For $\nu \in \sigma_\infty$, let n_ν denote the multiplicity of $\lambda_\nu = \nu^2 - \frac{(n-2)^2}{4}$ as the eigenvalue of $-\Delta_{\mathbb{S}^{n-1}} + q(\theta)$. For the class of potentials with critical decay, if zero is a resonance of P , its multiplicity is at most $\sum_{\nu \in \sigma_1} n_\nu$. Let u be a resonant state of P . Then one has the following asymptotics for $r = |x|$ large

$$u(x) = \frac{\psi(\theta)}{r^{\frac{n-2}{2} + \nu}} (1 + o(1)) \quad (1.8)$$

for some $\nu \in \sigma_1$ and $\psi \neq 0$ an eigenfunction of $-\Delta_{\mathbb{S}^{n-1}} + q(\theta)$ with eigenvalue λ_ν . We call u a ν -resonant state (or a ν -bound state). This terminology is consistent with the historical half-bound state for rapidly decreasing potentials in three dimensional case, since the set σ_1 is then reduced to $\{\frac{1}{2}\}$. The multiplicity m_ν of ν -resonant states is defined as the dimension of the subspace spanned by all ψ such that the expansion (1.8) holds for some resonant state u . The main result of this paper is the following generalized Levinson's theorem in presence of zero eigenvalue and zero resonance.

Theorem 1.1. *Let $n \geq 2$. Assume that v is smooth and satisfies (1.5) for some $\rho_0 > \max\{6, n+2\}$. Let $\xi(\cdot)$ denote the spectral shift function for the pair (P, P_0) . Then there exist some constants c_j and $\beta_{n/2}$ such that*

$$\int_0^\infty (\xi'(\lambda) - \sum_{j=1}^{\lfloor \frac{n}{2} \rfloor} c_j \lambda^{\lfloor \frac{n}{2} \rfloor - j - 1}) d\lambda = -(\mathcal{N}_- + \mathcal{N}_0 + \sum_{\nu \in \sigma_1} \nu m_\nu) + \beta_{n/2}. \quad (1.9)$$

Here $\beta_{n/2} = 0$ if n is odd and $c_{\lfloor \frac{n}{2} \rfloor} = 0$ if n is even. β_j for $j = 1, 2$ is computed at the end of paper. \mathcal{N}_- is the number of negative eigenvalues of P , \mathcal{N}_0 the multiplicity of the zero eigenvalue and m_ν that of ν -resonance of P . It is understood that if 0 is not an eigenvalue (resp., if there is no ν -resonant states), then $\mathcal{N}_0 = 0$ (resp. $m_\nu = 0$).

For the class of potentials under consideration, zero resonance of P may appear in any space dimension with arbitrary multiplicity depending on q . See [6, 32] in the case of manifolds with conical end and [14] in the case of \mathbb{R}^n . In [14], the existence of zero resonance is studied. Levinson's theorem in presence of zero eigenvalue and zero

resonance is mostly proved for spherically symmetric potentials by using Sturm-Liouville method or techniques of Jost function. See [20, 23] for overviews on the topic, where the reader can also find some discussions on radial potentials with critical decay. The proof of Levinson's theorem for non-radial potentials is much more evolved. See [22] and a series of works of D. Bollé *et al.* [1, 4, 5] for potentials satisfying $\langle x \rangle^s v \in L^1(\mathbb{R}^n)$ for some suitably large s depending on n and for $n = 1, 2, 3$. For example, one needs $s > 4$ if $n = 3$, $s > 8$ if $n = 2$ according to [5]. A formula as (1.9) in its general setting seems to be new.

To prove Theorem 1.1, we use the result of [25] on the high energy asymptotics of $\xi'(\lambda)$. The main task of our work is to analyze the spectral shift function in a neighbourhood of 0. Our approach is to calculate the trace of the operator-valued function $z \rightarrow (R(z) - R_0(z))f(P)$ and its generalized residue at 0, where $R_0(z) = (P_0 - z)^{-1}$, $R(z) = (P - z)^{-1}$ and $f \in C_0^\infty(\mathbb{R})$ is equal to 1 near 0. As the first step, we prove the asymptotics expansions of the resolvents $R(z)$ and $R_0(z)$ for z near 0 by adapting the methods of [33] to our situation. The main difficulties arise from the fact that the set σ_∞ may be arbitrary (depending on $q(\theta)$). The interplays between the zero energy resonant states and between the zero resonant states and the zero energy eigenfunctions can be produced in such a way that their contributions may be dominant over that of a single ν -resonant state with itself. This is overcome by carefully examining the terms obtained from asymptotic expansion of the resolvent and by making use of the properties of ν -resonant states. Remark that the decay condition can be improved if σ_1 contains only one point and our proofs (both for the resolvent expansions and for the Levinson's theorem) can be very much simplified if the potential $v(x)$ is spherically symmetric.

This work is organized as follows. In Section 2, we establish a representation formula of the spectral shift function which will be used to prove Levinson's theorem. In Section 3, we use the asymptotic expansion of $\tilde{R}_0(z) = (\tilde{P}_0 - z)^{-1}$ to get the asymptotic expansion for the resolvents $R_0(z)$ and $R(z)$. The methods are the same as in [33]. We indicate only the modifications to make and omit the details of calculation. The hard part of the proof of Theorem 1.1 is to calculate the generalized residue at 0 of the trace function $z \rightarrow \text{Tr}(R(z) - R_0(z))f(P)$, which is carried out in Section 4. Here f is some smooth function with compact support equal to 1 near 0. This result is used to study the low-energy asymptotics of the derivative of the spectral shift function in Section 5 and to prove the Levinson's theorem.

Notation. The scalar product on $L^2(\mathbb{R}_+; r^{n-1}dr)$ and $L^2(\mathbb{R}^n)$ is denoted by $\langle \cdot, \cdot \rangle$ and that on $L^2(\mathbb{S}^{n-1})$ by (\cdot, \cdot) . $H^{r,s}$, $r, s \in \mathbb{R}$, denotes the weighted Sobolev space of order r with the weight $\langle x \rangle^s$. The duality between $H^{1,s}$ and $H^{-1,-s}$ is identified with L^2 -product. Denote $H^{0,s} = L^{2,s}$. $\mathcal{L}(r, s; r', s')$ stands for the space of continuous linear operators from $H^{r,s}$ to $H^{r',s'}$. The complex plane \mathbb{C} is slit along positive real axis so that $z^\nu = e^{\nu \ln z}$ and $\ln z = \log|z| + i \arg z$ with $0 < \arg z < 2\pi$ are holomorphic for z near 0 in the slit complex plane.

2. A REPRESENTATION FORMULA

Let (P, P_0) be a pair of self-adjoint operators, semi-bounded from below, in some separable Hilbert space \mathcal{H} . We assume that for some $k \in \mathbb{N}^*$,

$$\|(P - i)^{-k} - (P_0 - i)^{-k}\|_{tr} < \infty, \quad (2.1)$$

where $\|\cdot\|_{tr}$ denotes the trace-class norm in \mathcal{H} . Under the assumption (2.1), it is well-known that for any $f \in \mathcal{S}(\mathbb{R})$, $f(P) - f(P_0)$ is of trace class.

For $z \in \mathbb{C}$ in the resolvent set of P_0 and P , we denote by $R_0(z) = (P_0 - z)^{-1}$, (resp. $R(z) = (P - z)^{-1}$) the resolvent of P_0 , (resp. of P). Writing

$$(R(z) - R_0(z))f(P) = [R(z)f(P) - R_0(z)f(P_0)] - R_0(z)(f(P) - f(P_0)), \quad (2.2)$$

we see immediately that $(R(z) - R_0(z))f(P)$ is of trace class.

The spectral shift function (SSF, in short) $\xi(\lambda) \in L^1_{loc}(\mathbb{R})$ is defined up to a constant as a Schwartz distribution on \mathbb{R} by the equation

$$\text{Tr} (f(P) - f(P_0)) = - \int_{\mathbb{R}} f'(\lambda) \xi(\lambda) d\lambda, \quad \forall f \in \mathcal{S}(\mathbb{R}). \quad (2.3)$$

The right hand side can be interpreted as $\langle \xi', f \rangle$, where ξ' is the derivative of ξ in the sense of the distributions, and $\langle \cdot, \cdot \rangle$ denotes the pairing between $\mathcal{S}'(\mathbb{R})$ and $\mathcal{S}(\mathbb{R})$.

In this section, we give a representation formula for the SSF under the following assumptions :

- The spectra of P and P_0 are purely absolutely continuous in $]0, +\infty[$:

$$\sigma(P_0) = \sigma_{ac}(P_0) = [0, +\infty[, \quad (2.4)$$

$$\sigma_{ac}(P) = [0, +\infty[. \quad (2.5)$$

In particular, we assume there are no embedded eigenvalues of P and P_0 in $]0, +\infty[$.

- The total number, \mathcal{N}_- , of negative eigenvalues of P is finite.
- For any $f \in C_0^\infty(\mathbb{R})$, there exists some $\epsilon_0 > 0$, $\delta_0 < 1$ and $C > 0$ (depending on f), such that

$$| \text{Tr} [(R(z) - R_0(z))f(P)] | \leq \frac{C}{|z|^{1+\epsilon_0}}, \quad (2.6)$$

for $z \in \mathbb{C}$ with $|z|$ large and $z \notin \sigma(P)$ and

$$| \text{Tr} [R_0(z)(f(P) - f(P_0))] | \leq \frac{C}{|z|^{\delta_0}}, \quad (2.7)$$

for $z \in \mathbb{C}$ with $|z|$ small and $z \notin \sigma(P_0)$.

- Let $f \in C_0^\infty(\mathbb{R})$ with $f(t) = 1$ for t near 0. The generalized residue of the function $z \rightarrow \text{Tr} [(R(z) - R_0(z))f(P)]$ at $z = 0$ is finite in the following sense: if we denote, for $\epsilon \ll \delta$,

$$\gamma(\delta, \epsilon) = \{z \in \mathbb{C} ; |z| = \delta, \text{dist}(z, \mathbb{R}^+) \geq \epsilon\},$$

we assume that

$$J_0 = -\frac{1}{2\pi i} \lim_{\delta \rightarrow 0} \lim_{\epsilon \rightarrow 0} \int_{\gamma(\delta, \epsilon)} \text{Tr} [(R(z) - R_0(z))f(P)] dz \text{ exists,} \quad (2.8)$$

where $\gamma(\delta, \epsilon)$ is positively oriented.

- The derivative of ξ in the sense of distributions satisfies

$$\xi'(\lambda) \in L^1_{loc}([0, \infty[). \quad (2.9)$$

Remark that conditions (2.1), (2.4)-(2.7) imply that if J_0 exists for some f as above, then it exists for all $f \in C_0^\infty(\mathbb{R})$ with $f(t) = 1$ for t near 0 and J_0 is independent of f . In fact, if f_1 and f_2 are two such functions, applying (2.2) to $f = f_1 - f_2$, one sees that $\text{Tr} [(R(z) - R_0(z))f(P)]$ can be decomposed into two terms, one is holomorphic in z near 0 and the other is of the order $O(\frac{1}{|z|^{\delta_0}})$, $\delta_0 < 1$. Hence the generalized residue of $z \rightarrow \text{Tr} [(R(z) - R_0(z))f(P)]$ at 0 is equal to zero. We have the following representation formula for the spectral shift function $\xi(\lambda)$.

Theorem 2.1. *Let $f \in C_0^\infty(\mathbb{R})$ such that $f(\lambda) = 1$ for λ in neighborhood of $\sigma_{pp}(P) \cup \{0\}$. Under the above assumptions, the limit*

$$\int_0^\infty f(\lambda) \xi'(\lambda) d\lambda = \lim_{\delta \rightarrow 0^+} \lim_{R \rightarrow \infty} \int_\delta^R f(\lambda) \xi'(\lambda) d\lambda$$

exists and one has

$$\text{Tr} (f(P) - f(P_0)) - \int_0^\infty \xi'(\lambda) f(\lambda) d\lambda = \mathcal{N}_- + J_0. \quad (2.10)$$

Proof. Let us consider the application $F(z) = \text{Tr} [(R(z) - R_0(z))f(P)]$ which is holomorphic outside $\sigma(P)$. We want to deduce (2.10) from Cauchy's formula applied to $F(z)$.

Let us begin by some notation. Let E_j , $1 \leq j \leq k$, be the distinct eigenvalues of P with multiplicity m_j . Denote for $z_0 \in \mathbb{C}$ and $\delta > 0$,

$$C(z_0; \delta) = \{ z \in \mathbb{C}; |z - z_0| = \delta \} \quad ; \quad D(z_0; \delta) = \{ z \in \mathbb{C}; |z - z_0| \leq \delta \}. \quad (2.11)$$

For $\delta > 0$ small enough, $\sigma(P) \cap D(E_i; \delta) = \{E_i\}$. For $R \gg 1$ and $0 < \epsilon \ll \delta$, let us denote

$$\gamma(R, \epsilon) = \{ z \in \mathbb{C}; |z| = R, \text{dist}(z, \mathbb{R}^+) \geq \epsilon \} \quad (2.12)$$

$$d^\pm(R, \delta, \epsilon) = [\sqrt{\delta^2 - \epsilon^2} \pm i\epsilon, \sqrt{R^2 - \epsilon^2} \pm i\epsilon]. \quad (2.13)$$

We denote by $\Gamma_{\delta, \epsilon, R}$ the curve defined by

$$\Gamma_{\delta, \epsilon, R} = \left(\bigcup_{j=1}^k C(E_j; \delta) \right) \cup \gamma(\delta, \epsilon) \cup d^+(R, \delta, \epsilon) \cup \gamma(R, \epsilon) \cup d^-(R, \delta, \epsilon).$$

$\Gamma_{\delta, \epsilon, R}$ is positively oriented according to the anti-clockwise orientation of the big circle $\gamma(R, \epsilon)$. Since $F(z)$ is holomorphic in the domain limited by $\Gamma_{\delta, \epsilon, R}$, the Cauchy integral formula gives

$$\frac{1}{2\pi i} \oint_{\Gamma_{\delta, \epsilon, R}} F(z) dz = 0.$$

We split the integral into four terms

$$\begin{aligned} \frac{1}{2i\pi} \oint_{\Gamma_{\delta,\epsilon,R}} F(z) dz &= \sum_{j=1}^4 I_j \quad \text{with} \\ I_1 &= \frac{1}{2i\pi} \oint_{\gamma(R,\epsilon)} F(z) dz, \quad I_2 = \sum_{j=1}^k \frac{1}{2i\pi} \oint_{C(E_j;\delta)} F(z) dz. \\ I_3 &= \frac{1}{2i\pi} \oint_{\gamma(\delta,\epsilon)} F(z) dz, \quad I_4 = \frac{1}{2i\pi} \oint_{d^+(R,\delta,\epsilon) \cup d^-(R,\delta,\epsilon)} F(z) dz. \end{aligned} \quad (2.14)$$

By condition (2.6), one has $I_1 = O(R^{-\epsilon_0})$ when $R \rightarrow \infty$, uniformly in δ and $\epsilon > 0$. For $j = 1, \dots, k$, E_j is an eigenvalue of P with finite multiplicity m_j , and the spectral projection associated to E_j is given by

$$\Pi_j = \frac{1}{2i\pi} \oint_{C(E_j;\delta)} R(z) dz, \quad (2.15)$$

since $C(E_j;\delta)$ is oriented according the clockwise orientation. Using that $R_0(z)$ is holomorphic in $D(E_j, \delta)$, we obtain :

$$I_2 = \sum_{j=1}^k \text{Tr} (\Pi_j f(P)) = \sum_{j=1}^k m_j f(E_j) = \sum_{j=1}^k m_j = \mathcal{N}_-. \quad (2.16)$$

Using assumption (2.8), we have by definition

$$\lim_{\delta \rightarrow 0} \lim_{\epsilon \rightarrow 0} I_3 = J_0, \quad (2.17)$$

and as a conclusion, we have obtained :

$$\lim_{R \rightarrow +\infty} \lim_{\delta \rightarrow 0} \lim_{\epsilon \rightarrow 0} I_4 = -\mathcal{N}_- - J_0. \quad (2.18)$$

Now, let us establish a new expression for I_4 ; we split $F(z) = F_1(z) + F_2(z)$ where

$$F_1(z) = \text{Tr} [R(z)f(P) - R_0(z)f(P_0)], \quad (2.19)$$

$$F_2(z) = \text{Tr} [R_0(z)(f(P_0) - f(P))]. \quad (2.20)$$

First, let us investigate the contribution coming from $F_1(z)$. By the definition of the SSF,

$$F_1(z) = - \int_{\mathbb{R}} \xi(\lambda) \frac{\partial f}{\partial \lambda}(\lambda, z) d\lambda \quad (2.21)$$

where $f(\lambda, z) = \frac{1}{\lambda - z} f(\lambda)$. Thus, $F_1(z) = G_1(z) - G_2(z)$ with

$$G_1(z) = \int_{\mathbb{R}} \xi(\lambda) \frac{1}{(\lambda - z)^2} f(\lambda) d\lambda, \quad (2.22)$$

$$G_2(z) = \int_{\mathbb{R}} \xi(\lambda) \frac{1}{\lambda - z} f'(\lambda) d\lambda. \quad (2.23)$$

For simplicity, we set $\Gamma = d^+(R, \delta, \epsilon) \cup d^-(R, \delta, \epsilon)$. We have

$$\frac{1}{2i\pi} \oint_{\Gamma} F_1(z) dz = \frac{1}{2i\pi} \oint_{\Gamma} G_1(z) dz - \frac{1}{2i\pi} \oint_{\Gamma} G_2(z) dz := (a) - (b). \quad (2.24)$$

Let us study (a). Using Fubini's theorem, we can write

$$(a) = \frac{1}{2i\pi} \int_{\mathbb{R}} \xi(\lambda) f(\lambda) \oint_{\Gamma} \frac{\partial}{\partial z} \left(\frac{1}{\lambda - z} \right) dz d\lambda. \quad (2.25)$$

Set

$$k(\lambda, s, \epsilon) = \frac{1}{\lambda - (\sqrt{s^2 - \epsilon^2} + i\epsilon)} - \frac{1}{\lambda - (\sqrt{s^2 - \epsilon^2} - i\epsilon)}.$$

One has

$$(a) = \frac{1}{2i\pi} \int_{\mathbb{R}} \xi(\lambda) f(\lambda) (k(\lambda, R, \epsilon) - k(\lambda, \delta, \epsilon)) d\lambda. \quad (2.26)$$

Let us look at the term I_ϵ related to $k(\lambda, R, \epsilon)$ and let us show that $\lim_{\epsilon \rightarrow 0} I_\epsilon = \xi(R)f(R)$.

Fix $0 < R_0 < R$ and split $I_\epsilon = J_\epsilon + K_\epsilon$, where

$$J_\epsilon = \frac{1}{2i\pi} \int_{-\infty}^{R_0} \xi(\lambda) f(\lambda) k(\lambda, R, \epsilon) d\lambda, \quad K_\epsilon = \frac{1}{2i\pi} \int_{R_0}^{+\infty} \xi(\lambda) f(\lambda) k(\lambda, R, \epsilon) d\lambda. \quad (2.27)$$

By the basic properties of the SSF, one has $\xi f \in L^1(\mathbb{R})$ and

$$|J_\epsilon| \leq \frac{1}{\pi} \int_{-\infty}^{R_0} |\xi(\lambda) f(\lambda)| \frac{\epsilon}{(\lambda - \sqrt{R^2 - \epsilon^2})^2 + \epsilon^2} d\lambda \leq C\epsilon \|\xi f\|_1. \quad (2.28)$$

On the other hand, we have

$$K_\epsilon = \frac{1}{\pi} \int_{\frac{R_0 - \sqrt{R^2 - \epsilon^2}}{\epsilon}}^{+\infty} (\xi f)(\sqrt{R^2 - \epsilon^2} + \epsilon s) \frac{1}{s^2 + 1} ds. \quad (2.29)$$

Under the assumption (2.9), $\xi \in C^0([0, +\infty[)$, so by making use of the dominated convergence theorem, one has

$$\lim_{\epsilon \rightarrow 0} K_\epsilon = \xi(R)f(R). \quad (2.30)$$

As a conclusion, we have shown

$$\lim_{\epsilon \rightarrow 0} (a) = \xi(R)f(R) - \xi(\delta)f(\delta) = -\xi(\delta), \quad (2.31)$$

for $\delta > 0$ sufficiently small and $R \gg 1$.

Now, let us study (b). Using Fubini's theorem,

$$(b) = \int_{\mathbb{R}} \xi(\lambda) f'(\lambda) \left(\frac{1}{2i\pi} \oint_{\Gamma} \frac{1}{\lambda - z} dz \right) d\lambda. \quad (2.32)$$

It is easy to check that $\frac{1}{2i\pi} \oint_{\Gamma} \frac{1}{\lambda - z} dz$ is uniformly bounded with respect to ϵ and

$$\lim_{\epsilon \rightarrow 0} \frac{1}{2i\pi} \oint_{\Gamma} \frac{1}{\lambda - z} dz = \mathbf{1}_{[\delta, R]}(\lambda) \text{ a.e. } \lambda. \quad (2.33)$$

So, making use of the dominated convergence theorem and then the assumption (2.9), a straightforward application of distribution theory gives

$$\lim_{\epsilon \rightarrow 0} (b) = \int_{\delta}^R \xi(\lambda) f'(\lambda) d\lambda = -\xi(\delta) - \int_{\delta}^R \xi'(\lambda) f(\lambda) d\lambda, \quad (2.34)$$

for $R \gg 1$ and $\delta > 0$ sufficiently small. As a conclusion, we have obtained,

$$\lim_{\epsilon \rightarrow 0} \frac{1}{2i\pi} \oint_{\Gamma} F_1(z) dz = \int_{\delta}^R \xi'(\lambda) f(\lambda) d\lambda. \quad (2.35)$$

To study the contribution coming from $F_2(z)$, we write

$$\frac{1}{2i\pi} \oint_{\Gamma} F_2(z) dz = \text{Tr} \left(\frac{1}{2i\pi} \oint_{\Gamma} R_0(z) dz (f(P_0) - f(P)) \right). \quad (2.36)$$

Using the spectral theorem for P_0 and the assumption (2.1), the same argument as (2.33) gives

$$\text{s-}\lim_{\epsilon \rightarrow 0} \frac{1}{2i\pi} \oint_{\Gamma} R_0(z) dz = \mathbf{1}_{[\delta, R]}(P_0). \quad (2.37)$$

Since $f(P) - f(P_0)$ is of trace class, and

$$\text{s-}\lim_{\delta \rightarrow 0, R \rightarrow \infty} \mathbf{1}_{[\delta, R]}(P_0) = Id, \quad (2.38)$$

one can deduce (see Lemma 2.2 below) that

$$\lim_{R \rightarrow \infty} \lim_{\delta \rightarrow 0} \lim_{\epsilon \rightarrow 0} \frac{1}{2i\pi} \oint_{\Gamma} F_2(z) dz = \text{Tr} (f(P_0) - f(P)). \quad (2.39)$$

So, using (2.18), (2.35) and (2.39), we obtain the result. \square

To complete the proof of Theorem 2.1, we prove the following elementary lemma.

Lemma 2.2. *Let A be an operator of trace class. Assume that $f(\lambda)$ is an operator valued function, uniformly bounded in λ , such that $\text{s-}\lim_{\lambda \rightarrow \lambda_0} f(\lambda) = B$ exists.*

Then, $f(\lambda)A$ converges to BA in the trace norm class, as $\lambda \rightarrow \lambda_0$. In particular,

$$\lim_{\lambda \rightarrow \lambda_0} \text{Tr} (f(\lambda)A) = \text{Tr}(BA). \quad (2.40)$$

Proof. For any $\epsilon > 0$, let F be a finite rank operator such that $\|A - F\|_{tr} < \epsilon$. Then,

$$\|f(\lambda)A - BA\|_{tr} \leq \|(f(\lambda) - B)F\|_{tr} + \|(f(\lambda) - B)(A - F)\|_{tr}.$$

Since F is a finite rank operator and $\text{s-}\lim_{\lambda \rightarrow \lambda_0} f(\lambda) = B$, we have $\|(f(\lambda) - B)F\|_{tr} \leq C\epsilon$ for $|\lambda - \lambda_0| \leq \delta$ with some $\delta > 0$ small enough. Since $f(\lambda) - B$ is a uniformly bounded in λ , we have also $\|(f(\lambda) - B)(A - F)\|_{tr} \leq C\epsilon$. This ends the proof. \square

The remaining part of this work is to apply Theorem 2.1 to Schrödinger operator, using the known results on the asymptotic expansion of $\xi'(\lambda)$ as $\lambda \rightarrow \infty$. The main task is to study $\xi(\lambda)$ for λ near 0. If one can calculate the generalized residue J_0 and can show that $\xi'(\lambda)$ is integrable in $]0, 1]$, then one can take a family of functions $f_R(\lambda) = \chi(\frac{\lambda}{R})$, where χ is smooth and $0 \leq \chi(s) \leq 1$, $\chi(s) = 1$ for s near 0, $\chi(s) = 0$ for $s > 1$ and expand both the terms

$$\int_0^\infty \xi'(\lambda) f_R(\lambda) d\lambda, \text{ and } \text{Tr}(f_R(P) - f_R(P_0))$$

in R large. Theorem 1.1 can be derived from Theorem 2.1 by comparing the two asymptotic expansions in R .

3. RESOLVENT ASYMPTOTICS NEAR THE THRESHOLD

Consider the Schrödinger operator $P = -\Delta + v(x)$ where the potential $v(x)$ is bounded and has the asymptotic behavior

$$v(x) = \frac{q(\theta)}{r^2} + O(\langle x \rangle^{-\rho_0}), \quad |x| \rightarrow \infty \quad (3.1)$$

where $\rho_0 > 2$, (r, θ) is the polar coordinates on \mathbb{R}^n and $q(\cdot)$ is continuous on the sphere. Let $\tilde{P}_0 = -\Delta + \frac{q(\theta)}{r^2}$ be the homogeneous part of P . Assume $n \geq 2$ and

$$-\Delta_{\mathbb{S}^{n-1}} + q(\theta) > -\frac{1}{4}(n-2)^2, \quad \text{on } L^2(\mathbb{S}^{n-1}). \quad (3.2)$$

In particular, (3.2) implies that the quadratic form defined by \tilde{P}_0 on $C_0^\infty(\mathbb{R}^n \setminus \{0\})$ is positive. We still denote by \tilde{P}_0 its Friedrich's realization as selfadjoint operator in $L^2(\mathbb{R}^n)$. Due to troubles related to the critical local singularity in the definition of spectral shift function, the model operator P_0 we will use in the next Section is a modification of \tilde{P}_0 inside some compact set. Let $0 \leq \chi_j \leq 1$ ($j = 1, 2$) be smooth functions on \mathbb{R}^n such that $\text{supp} \chi_1 \subset B(0, R_1)$, $\chi_1(x) = 1$ when $|x| < R_0$ and

$$\chi_1(x)^2 + \chi_2(x)^2 = 1.$$

Define

$$P_0 = \chi_1(-\Delta)\chi_1 + \chi_2\tilde{P}_0\chi_2,$$

on $L^2(\mathbb{R}^n)$ (if $q = 0$, we take $P_0 = -\Delta$). In the next Section, P will be regarded as a perturbation of P_0 . But to study the low energy asymptotics of $R_0(z)$ and $R(z)$, we regard both P and P_0 as perturbations of \tilde{P}_0 . Denote

$$P_0 = \tilde{P}_0 - W, \quad (3.3)$$

$$P = P_0 + V = \tilde{P}_0 - \tilde{W} \quad (3.4)$$

where

$$\begin{aligned} W &= \chi_1^2 \frac{q(\theta)}{r^2} - \sum_{j=1}^2 |\nabla \chi_j|^2 \\ V &= -\sum_{i=1}^2 |\nabla \chi_i|^2 + v - \chi_2^2 \frac{q(\theta)}{r^2} \\ \tilde{W} &= W - V = -v + \frac{q(\theta)}{r^2}. \end{aligned}$$

Note that W and \tilde{W} contain a critical singularity at zero. For $n \geq 3$, the Hardy's inequality implies that they map continuously H^1 to H^{-1} . When $n = 2$, the same remains true under the condition (1.4) if one defines H^1 as the form domain of \tilde{P}_0 . V is bounded and satisfies

$$|V(x)| \leq C \langle x \rangle^{-\rho_0}, \quad \rho_0 > 2. \quad (3.5)$$

P_0 is still a Schrödinger operator with a potential of critical decay at the infinity. But it has nice threshold spectral property at zero.

Let

$$\sigma_\infty = \left\{ \nu; \nu = \sqrt{\lambda + \frac{(n-2)^2}{4}}, \lambda \in \sigma(-\Delta_{\mathbb{S}^{n-1}} + q(\theta)) \right\}. \quad (3.6)$$

Denote

$$\sigma_k = \sigma_\infty \cap]0, k], \quad k \in \mathbb{N}. \quad (3.7)$$

For $\nu \in \sigma_\infty$, let n_ν denote the multiplicity of $\lambda_\nu = \nu^2 - \frac{(n-2)^2}{4}$ as the eigenvalue of $-\Delta_{\mathbb{S}^{n-1}} + q(\theta)$. Let $\varphi_\nu^{(j)}$, $\nu \in \sigma_\infty$, $1 \leq j \leq n_\nu$ denote an orthonormal basis of $L^2(\mathbb{S}^{n-1})$ consisting of eigenfunctions of $-\Delta_{\mathbb{S}^{n-1}} + q(\theta)$:

$$(-\Delta_{\mathbb{S}^{n-1}} + q(\theta))\varphi_\nu^{(j)} = \lambda_\nu \varphi_\nu^{(j)}, \quad (\varphi_\nu^{(i)}, \varphi_\nu^{(j)}) = \delta_{ij}.$$

Let π_ν denote the orthogonal projection in $L^2(\mathbb{R}_+ \times \mathbb{S}^{n-1}; r^{n-1} dr d\theta)$ onto the subspace spanned by the eigenfunctions of $-\Delta_{\mathbb{S}^{n-1}} + q(\theta)$ associated with the eigenvalue λ_ν :

$$\pi_\nu f = \sum_{j=1}^{n_\nu} (f, \varphi_\nu^{(j)}) \otimes \varphi_\nu^{(j)}, \quad f \in L^2(\mathbb{R}_+ \times \mathbb{S}^{n-1}; r^{n-1} dr d\theta).$$

Let

$$Q_\nu = -\frac{d^2}{dr^2} - \frac{n-1}{r} \frac{d}{dr} + \frac{\nu^2 - \frac{(n-2)^2}{4}}{r^2}, \quad \text{in } L^2(\mathbb{R}_+; r^{n-1} dr).$$

The asymptotic expansion of the resolvent $R_0(z) = (P_0 - z)^{-1}$ near zero can be obtained as in [33], by considering P_0 as a perturbation of \tilde{P}_0 . For the later purpose, we also need the asymptotic expansion of $\frac{d}{dz} \tilde{R}_0(z)$ for z near 0, where $\tilde{R}_0(z) = (\tilde{P}_0 - z)^{-1}$ for $z \notin \sigma(P_0)$.

Decomposition the resolvent $\tilde{R}_0(z)$ as

$$\tilde{R}_0(z) = \sum_{\nu \in \sigma_\infty} (Q_\nu - z)^{-1} \pi_\nu, \quad z \notin \mathbb{R}$$

The Schwartz kernel $K_\nu(r, \tau; z)$ of $(Q_\nu - z)^{-1}$, $\Im z > 0$, can be calculated explicitly. In fact, the Schwartz kernel of e^{-itQ_ν} is given by (see [29])

$$\frac{1}{2it} (r\tau)^{-\frac{n-2}{2}} e^{-\frac{r^2+\tau^2}{4it} - i\frac{\pi\nu}{2}} J_\nu\left(\frac{r\tau}{2t}\right), \quad t \in \mathbb{R}, \quad (3.8)$$

where $J_\nu(\cdot)$ is the Bessel function of the first kind of order ν . Since

$$(Q_\nu - z)^{-1} = i \int_0^\infty e^{-it(Q_\nu - z)} dt$$

for $\Im z > 0$, the Schwartz kernel of $(Q_\nu - z)^{-1}$ is

$$\begin{aligned} K_\nu(r, \tau; z) &= (r\tau)^{-\frac{n-2}{2}} \int_0^\infty e^{-\frac{r^2+\tau^2}{4it} + izt - i\frac{\pi\nu}{2}} J_\nu\left(\frac{r\tau}{2t}\right) \frac{dt}{2t} \\ &= (r\tau)^{-\frac{n-2}{2}} \int_0^\infty e^{\frac{i\rho}{t} + izr\tau t - i\frac{\pi\nu}{2}} J_\nu\left(\frac{1}{2t}\right) \frac{dt}{2t} \end{aligned} \quad (3.9)$$

for $\Im z > 0$, where

$$\rho = \rho(r, \tau) \equiv \frac{r^2 + \tau^2}{4r\tau}.$$

Note that the formula of $K_\nu(r, \tau; z)$ used in [6, 32] contains a sign error.

The asymptotic expansion for $(Q_\nu - z)^{-1}$ as $z \rightarrow 0$ and $\Im z > 0$ is deduced from (3.9) by splitting the last integral into two parts according to $t \in]0, 1]$ or $t \in [1, \infty[$. The integral for $t \in]0, 1]$ gives rise to a formal power series expansion in z near 0. The expansion corresponding to the integral of $t \in [1, \infty[$ needs a lengthy calculation. See [32]. Set

$$f(s; r, \tau, \nu) = D_\nu(r, \tau) \int_{-1}^1 e^{i(\rho+\theta/2)s} (1-\theta^2)^{\nu-1/2} d\theta, \quad \nu \geq 0,$$

with

$$D_\nu(r, \tau) = a_\nu (r\tau)^{-\frac{(n-2)}{2}}, \quad a_\nu = \frac{e^{-i\pi\nu/2}}{2^{2\nu+1} \pi^{1/2} \Gamma(\nu + 1/2)}. \quad (3.10)$$

Then f can be expanded in a convergent power series in s

$$f(s; r, \tau, \nu) = \sum_{j=0}^{\infty} s^j f_j(r, \tau, \nu), \quad s \in \mathbb{R}, \quad (3.11)$$

with

$$f_j(r, \tau, \nu) = (r\tau)^{-\frac{1}{2}(n-2)} P_{j,\nu}(\rho), \quad (3.12)$$

with $P_{j,\nu}(\rho)$ a polynomial in ρ of degree j :

$$P_{j,\nu}(\rho) = \frac{i^j a_\nu}{j!} \int_{-1}^1 (\rho + \frac{1}{2}\theta)^j (1 - \theta^2)^{\nu - \frac{1}{2}} d\theta.$$

In particular,

$$\begin{aligned} f_0(r, \tau, \nu) &= d_\nu (r\tau)^{-\frac{n-2}{2}}, \quad d_\nu = \frac{e^{-\frac{1}{2}i\pi\nu}}{2^{2\nu+1}\Gamma(\nu+1)}; \\ f_1(r, \tau, \nu) &= id_\nu (r\tau)^{-\frac{n-2}{2}} \rho. \end{aligned}$$

In [32, 33], the expansion of $(\tilde{P}_0 - z)^{-1}$ is given for z near 0. For the later use, we need the derivation of this expansion. By an abuse of notation, we denote by the same letter F an operator on $L^2(\mathbb{R}_+, r^{n-1}dr)$ and its distributional kernel.

For $\nu \in \sigma_\infty$, denote $[\nu]$ the integral part of ν and $\nu' = \nu - [\nu]$. Set

$$z_\nu = \begin{cases} z^{\nu'}, & \text{if } \nu \notin \mathbb{N}, \\ z \ln z, & \text{if } \nu \in \mathbb{N}. \end{cases}$$

For $\nu > 0$, let $[\nu]_-$ be the largest integer strictly less than ν . When $\nu = 0$, set $[\nu]_- = 0$. Define δ_ν by $\delta_\nu = 1$, if $\nu \in \sigma_\infty \cap \mathbb{N}$, $\delta_\nu = 0$, otherwise. One has $[\nu] = [\nu]_- + \delta_\nu$.

Proposition 3.1. *Assume the condition (3.2). One has*

(a). *The following asymptotic expansion holds for z near 0 with $\Im z > 0$.*

$$\tilde{R}_0(z) = \sum_{j=0}^N z^j F_j + \sum_{\nu \in \sigma_N} \sum_{j=[\nu]_-}^{N-1} z_\nu z^j G_{\nu,j+\delta_\nu} \pi_\nu + \tilde{R}_0^{(N)}(z), \quad (3.13)$$

in $\mathcal{L}(-1, s; 1, -s)$, $s > 2N + 1$. The remainder term $R_0^{(N)}(z)$ can be estimated by

$$\tilde{R}_0^{(N)}(z) = O(|z|^{N+\epsilon}) \in \mathcal{L}(-1, s; 1, -s), \quad s > 2N + 1,$$

for some $\epsilon > 0$ and F_j is of the form

$$F_j = \sum_{\nu \in \sigma_\infty} F_{\nu,j} \pi_\nu \in \mathcal{L}(-1, s; 1, -s), \quad s > 2j + 1. \quad (3.14)$$

$F_{\nu,j}$ and $G_{\nu,j}$ can be explicitly calculated from the asymptotic expansion of $K_\nu(r, \tau, z)$ in z . In particular,

$$G_{\nu,j}(r, \tau) = \begin{cases} b_{\nu',j}(r\tau)^{j+\nu'} f_{j-[\nu]}(r, \tau; \nu), & \nu \notin \mathbb{N} \\ -\frac{(ir\tau)^j}{j!} f_{j-[\nu]}(r, \tau; \nu), & \nu \in \mathbb{N} \end{cases}$$

with f_k defined by (3.12) and

$$b_{\nu',j} = -\frac{i^j e^{-i\nu'\pi/2} \Gamma(1 - \nu')}{\nu'(\nu' + 1) \cdots (\nu' + j)}, \quad (3.15)$$

for $0 < \nu' < 1$.

(b). *The expansion (3.13) can be differentiated with respect to z and one has*

$$\frac{d}{dz} \tilde{R}_0^{(N)}(z) = O(|z|^{N-1+\epsilon}) \in \mathcal{L}(-1, s; 1, -s), \quad s > 2N + 1,$$

for some $\epsilon > 0$ small enough.

See [32] for the proof of (a). To show that it is possible to differentiate the asymptotic expansion in z , one first shows that it can be done for each $(Q_\nu - z)^{-1}$, then one utilizes the properties of Bessel function to control the dependence on ν , including the remainders, and finally takes the sum in ν . See [14] for details.

Remark 3.2. For the later use, let us precise a few terms in $R_0(z)$. By an abuse of notation, we denote by same letter an operator on $L^2(\mathbb{R}_+; r^{n-1}dr)$ and its distribution kernel. Then one has

$$F_{\nu,0} = (r\tau)^{-\frac{1}{2}(n-2)} \int_0^\infty e^{i\frac{t}{2} - i\frac{\pi\nu}{2}} J_\nu\left(\frac{1}{2t}\right) \frac{dt}{2t}, \quad \nu \in \sigma_\infty, \quad (3.16)$$

$$G_{\nu,0} = -\frac{e^{-i\pi\nu}\Gamma(1-\nu)}{\nu 2^{2\nu+1}\Gamma(1+\nu)} (r\tau)^{-\frac{n-2}{2}+\nu}, \quad \nu \in]0, 1[, \quad (3.17)$$

$$G_{1,1} = -\frac{1}{8}(r\tau)^{-\frac{n-2}{2}+1}. \quad (3.18)$$

By (3.16), one can derive that

$$|F_{\nu,0}(r, \tau)| \leq C(r\tau)^{-\frac{n-2}{2}} \left(\frac{r\tau}{r^2 + \tau^2}\right)^{\min\{1, \nu\}}, \quad (3.19)$$

for some $C > 0$ independent of $\nu \in \sigma_\infty$. The uniformity in ν is obtained by examining the dependence of $J_\nu(r)$ on ν .

The asymptotic expansion for $R_0(z) = (P_0 - z)^{-1}$ can be deduced by P_0 as perturbations of \tilde{P}_0 . One has

$$R_0(z) = (1 - F(z))^{-1} \tilde{R}_0(z), \quad R(z) = (1 - \tilde{F}(z))^{-1} \tilde{R}_0(z) \quad (3.20)$$

where

$$F(z) = \tilde{R}_0(z)W, \quad \tilde{F}(z) = \tilde{R}_0(z)\tilde{W}. \quad (3.21)$$

For $n \geq 3$, the multiplication by $\frac{1}{|x|^2}$ belongs to $\mathcal{L}(1, s; -1, s)$ for any s , by the Hardy inequality. By (3.26), the same is true for $n = 2$ if we define $H^{1,s}$ as $\langle x \rangle^{-s} Q(\tilde{P}_0)$, where $Q(\tilde{P}_0)$ is the form-domain of \tilde{P}_0 . Therefore although W and \tilde{W} have a critical singularity $\frac{1}{|x|^2}$ at zero, Proposition 3.1 implies that $F(z) = F_0W + O(|z|^\epsilon)$ in $\mathcal{L}(1, -s; 1, -s)$ for $s > 1$ (here and in the following, $H^{1,s}$ is replaced by $\langle x \rangle^{-s} Q(\tilde{P}_0)$ when $n = 2$). Similar result holds for $\tilde{F}(z)$. Note that F_0W and $F_0\tilde{W}$ are not compact operators.

Definition 3.3. Set $\mathcal{N}(P) = \{ u; F_0\tilde{W}u = u, u \in H^{1,-s}, \forall s > 1 \}$. A function $u \in \mathcal{N}(P) \setminus L^2$ is called a resonant state of P at zero. If $\mathcal{N}(P) = \{0\}$, we say that 0 is the regular point of P . The multiplicity of the zero resonance of P is defined as $\mu_r = \dim \mathcal{N} / (\ker_{L^2} P)$. Zero resonance and resonant states of P_0 are defined in the same way with \tilde{W} replaced by W .

For $u \in H^{1,-s}$ for any $s > 1$ and $u \in \mathcal{N}$, one can show that $Pu = (\tilde{P}_0 - \tilde{W})u = 0$. If $\tilde{W} = O(\langle x \rangle^{-\rho_0})$ with $\rho_0 > 3$, it is proved in [32] that

$$u(r\theta) = \sum_{\nu \in \sigma_1} \sum_{j=1}^{n_\nu} \frac{1}{2\nu} \langle \tilde{W}u, |y|^{-\frac{n-2}{2}+\nu} \varphi_\nu^{(j)} \rangle \frac{\varphi_\nu^{(j)}(\theta)}{r^{\frac{n-2}{2}+\nu}} + \tilde{u}, \quad (3.22)$$

where $\tilde{u} \in L^2(|x| > 1)$, and (\cdot, \cdot) is the scalar product in $L^2(\mathbb{S}^{n-1})$. In particular, (3.22) shows that the multiplicity of the zero resonance of P is bounded by the total multiplicity of eigenvalues λ_ν of $-\Delta_{\mathbb{S}^{n-1}} + q(\theta)$ with $\nu \in \sigma_1$.

For $\nu \in \sigma_1$, we shall say u is a ν -resonant state, or ν -bound state, of P if $u \in \mathcal{N}(P)$ and if u has an asymptotic behavior like

$$u(x) = \frac{\psi(\theta)}{r^{\frac{n-2}{2}+\nu}} + o\left(\frac{1}{r^{\frac{n-2}{2}+\nu}}\right),$$

for some $\psi \neq 0$, as $r \rightarrow \infty$. In the case $n = 3$ and $q(\theta) = 0$, one has $\sigma_1 = \{\frac{1}{2}\}$. The only possible zero energy resonant states of P are half-bound states, which is in agreement with the usual terminology on this topic. In the general case, (3.22) shows that ψ is an eigenfunction of $-\Delta_{\mathbb{S}^{n-1}} + q(\theta)$ associated with the eigenvalue λ_ν . We shall say that m ν -resonant states of P , denoted as u_1, \dots, u_m , are linearly independent if

$$u_l(x) = \frac{\psi_l(\theta)}{r^{\frac{n-2}{2}+\nu}}(1 + o(1)), \quad r \rightarrow \infty,$$

with $\{\psi_1, \dots, \psi_m\}$ linearly independent in $L^2(\mathbb{S}^{n-1})$. Let m_ν denote the maximal number of linearly independent ν -resonant states of P . Then, m_ν does not exceed the multiplicity of the eigenvalue λ_ν of $-\Delta_{\mathbb{S}^{n-1}} + q(\theta)$ and

$$\sum_{\nu \in \sigma_1} m_\nu = \mu_r. \quad (3.23)$$

Note in particular that if u is a ν -resonant state, then one has

$$\langle \widetilde{W}u, |y|^{-\frac{n-2}{2}+\mu} \varphi_\mu^{(j)} \rangle = 0 \quad (3.24)$$

for all $\mu \in \sigma_1$ with $\mu < \nu$ and for all j with $1 \leq j \leq n_\mu$ and if u is an eigenfunction of P associated with the eigenvalue 0, the above equality remains true for all $\nu \in \sigma_1$ and all j with $1 \leq j \leq n_\mu$ (see (3.22)). These properties will be repeatedly used in the calculation of the generalized residue in Section 4. Finally, remark that if $m_\nu \neq 0$ for some $\nu \in \sigma_1$, we can choose m_ν ν -resonant states $u_l = \frac{\psi_l(\theta)}{r^{\frac{n-2}{2}+\nu}}(1 + o(1))$, $1 \leq l \leq m_\nu$, such that $\{\psi_l\}$ is orthonormal in $L^2(\mathbb{S}^{n-1})$. Modifying the basis $\{\varphi_\nu^{(j)}\}$ used in the definition of the spectral projection π_ν , we can assume without loss that $\varphi_\nu^{(l)} = \psi_l$, $1 \leq l \leq m_\nu$.

The model operator P_0 to be used in the next Section has the following nice threshold spectral property.

Lemma 3.4. *Assume $n \geq 2$ and (3.2). Zero is a regular point of P_0 .*

Proof. Recall that $P_0 = \widetilde{P}_0 - W$ with W of compact support. Let $u \in \mathcal{N}(P_0)$. Let $\nu_0 = \min \sigma_1 > 0$, by (3.2). Since W is of compact support, $Wu \in H^{-1,t}$ for any $t > 0$. It follows from (3.22) that $u \in H^{1,-s'}$ for any $s' > 1 - \nu_0$. On the other hand, $P_0 u = 0$ implies that $-\Delta u = (W - \frac{q(\theta)}{|x|^2})u \in L^{2,2-s'}$. By the assumption (3.2), one can deduce that there exists $c > 0$ such that

$$\langle |x|^{-2} \chi_2 u, \chi_2 u \rangle \leq c \langle \widetilde{P}_0 \chi_2 u, \chi_2 u \rangle. \quad (3.25)$$

In fact, let $\rho \in C_0^\infty(\mathbb{R}^n)$ with $\rho(x) = 1$ for $|x| \leq 1$. Set $u_m = \rho(x/m) \chi_2 u$, $m \in \mathbb{N}^*$. Then $u_m \in H^1$, and by the ellipticity of $-\Delta$ one has in fact $u_m \in H^2$. The assumption (3.2) implies that there exists $\epsilon_0 > 0$ such that for $f \in H^1$ with compact support in $\mathbb{R}^n \setminus \{0\}$, one has

$$\langle \widetilde{P}f, f \rangle \geq \int \int \left(\left| \frac{\partial f}{\partial r} \right|^2 + \left(\epsilon_0 - \frac{(n-2)^2}{4} \right) \frac{|f|^2}{r^2} \right) r^{n-1} dr d\theta.$$

Making use of the Hardy inequality, one obtains for $n \geq 2$

$$\langle |x|^{-2}f, f \rangle \leq \epsilon_0^{-1} \langle \tilde{P}_0 f, f \rangle. \quad (3.26)$$

Since $u \in H^{1,-s}$ and $-\Delta u \in L^{2,2-s}$ for any $s > 1 - \nu_0$, $\nu_0 > 0$, we can take $s \in]1 - \nu_0, 1[$. Applying (3.26) to $f = u_m$ and taking the limit $m \rightarrow \infty$, we derive (3.25) by noticing that the term related to $[-\Delta, \rho(x/m)]$ tends to 0, due to the decay of u . (3.25) implies in particular that $\langle \tilde{P}_0 \chi_2 u, \chi_2 u \rangle \geq 0$. The equation

$$\langle -\Delta \chi_1 u, \chi_1 u \rangle + \langle \tilde{P}_0 \chi_2 u, \chi_2 u \rangle = \langle P_0 u, u \rangle = 0$$

shows that each term of the above sum vanishes. The estimate (3.25) gives in turn $\chi_2 u = 0$. Now the unique continuation theorem shows that $u = 0$. This proves that zero is a regular point of P_0 . \square

The existence of the asymptotic expansion of the resolvent $R_0(z)$ can easily be obtained by a method of perturbation. Concretely, let $K(z)$ be defined by

$$K(z) = (\chi_1(-\Delta + 1 - z)^{-1} \chi_1 + \chi_2(\tilde{P}_0 - z)^{-1} \chi_2)(P_0 - z) - 1, \quad z \notin \mathbb{R}_+. \quad (3.27)$$

$K(z)$ is compact operator on $H^{1,-s}$, $s > 1$. By Proposition 3.1, the limit $K(0) = \lim_{z \rightarrow 0} K(z)$ exists and is compact. The kernel of $1 + K(0)$ in $H^{1,-s}$ coincides with $\mathcal{N}(P_0)$ (both are equal to solutions to the equation $P_0 u = 0$, $u \in H^{1,-s}$) which by Lemma 3.4 is $\{0\}$. So $(1 + K(0))^{-1}$ is invertible in $H^{1,-s}$. From Proposition 3.1, we can derive the asymptotic expansion of $R_0(z)$ in suitable weighted spaces from the formula

$$R_0(z) = (1 + K(z))^{-1} (\chi_1(-\Delta + 1 - z)^{-1} \chi_1 + \chi_2(\tilde{P}_0 - z)^{-1} \chi_2). \quad (3.28)$$

In particular, $R_0(0) = \lim_{z \rightarrow 0, z \notin \mathbb{R}_+} R_0(z)$ exists in $\mathcal{L}(-1, s; 1, -s)$, $s > 1$ and

$$R_0(0) = (1 + K(0))^{-1} (\chi_1(-\Delta + 1)^{-1} \chi_1 + \chi_2 F_0 \chi_2). \quad (3.29)$$

However since $K(z)$ contains several terms induced by cut-offs, the expansion obtained in this way is too complicated to be useful in the proof of Levinson's theorem which requires detailed information on higher order terms. For this purpose, we use the resolvent equation $R_0(z) = (1 - \tilde{R}_0(z)W)^{-1} \tilde{R}_0(z)$ to obtain a more concise expansion.

Proposition 3.5. *Assume (3.2), $n \geq 2$ and $\rho_0 > 2$.*

- (a). $1 - F_0 W$ is invertible on $H^{1,-s}$, $s > 1$.
- (b). Let $s \in]1, \rho_0/2[$. $1 - F_0 \tilde{W}$ is a Fredholm operator on $H^{1,-s}$ with indices (m, m) , $m = \dim \mathcal{N}(P) < \infty$. One has $H^{1,-s} = \ker(1 - F_0 \tilde{W}) \oplus \text{ran}(1 - F_0 \tilde{W})$.

Proof. (a). Lemma 3.4 shows that $1 - F_0 W$ is injective. Since $(1 - F_0 W)^* = 1 - W F_0$, F_0 is injective from $\text{coker}(1 - F_0 W)$ to $\ker(1 - F_0 W)$. Therefore $\text{coker}(1 - F_0 W) = \{0\}$ and $\text{ran}(1 - W F_0)$ is dense in $H^{1,-s}$. For any $u \in H^{1,-s}$, set

$$f = u + R_0(0)(Wu)$$

with $R_0(0)$ given by (3.29). Then $f \in H^{1,-s}$. Since F_0 and $R_0(0)$ are limits of $\tilde{R}_0(z)$ and $R_0(z)$ in $\mathcal{L}(-1, s; 1, -s)$, we can check that

$$(1 - F_0 W)f = u + (R_0(0) - F_0 - F_0 W R_0(0))Wu = u. \quad (3.30)$$

This proves that $u \in \text{ran}(1 - F_0 W)$ and $1 - F_0 W$ is bijective on $H^{1,-s}$. The open mapping theorem shows that $1 - F_0 W$ is invertible on $H^{1,-s}$.

To show (b), recall that $1 + R_0(0)V$ is a Fredholm operator with equal indices (m, m) , $m = \dim \ker(1 + R_0(0)V)$, and that $H^{1,-s} = \ker(1 + R_0(0)V) \oplus \text{ran}(1 + R_0(0)V)$ (see

[13, 30]). By the relation $P = P_0 + V = \tilde{P}_0 - \tilde{W}$, one can show that $\mathcal{N}(P) = \ker P = \ker(1 + R_0(0)V)$ in $H^{1,-s}$, $s \in]1, \rho_0/2[$ and that $\text{coker}(1 - F_0\tilde{W})$ is also of dimension m . To show that the range of $1 - F_0\tilde{W}$ is closed, let $u \in H^{1,-s}$ and $\{u_n\} \subset \text{ran}(1 - F_0\tilde{W})$ such that $u_n \rightarrow u$ in $H^{1,-s}$. Let $v_n \in H^{1,-s}$ such that $u_n = (1 - F_0\tilde{W})v_n$. Then $\tilde{P}_0 u_n = P v_n = (P_0 + V)v_n$. It follows that

$$(1 + R_0(0)W)u_n = (1 + R_0(0)V)v_n. \quad (3.31)$$

Since $1 + R_0(0)W$ is continuous on $H^{1,-s}$, the left-hand side of (3.31) converges to $(1 + R_0(0)W)u$ as $n \rightarrow \infty$, while the right-hand side is clearly in the range of $1 + R_0(0)V$. Since $R_0(0)V$ is compact, the range of $1 + R_0(0)V$ is closed. It follows from (3.31) that there exists $v \in H^{1,-s}$ such that $(1 + R_0(0)W)u = (1 + R_0(0)V)v$. One can check that $u = (1 - F_0\tilde{W})v \in \text{ran}(1 - F_0\tilde{W})$, which proves that the range of $(1 - F_0\tilde{W})$ is closed. It follows that $(1 - F_0\tilde{W})$ is a Fredholm operator with equal indices. The other affirmation of (b) can be proved in the same way as in [30]. \square

Denote

$$\begin{aligned} \vec{\nu} &= (\nu_1, \dots, \nu_k) \in (\sigma_N)^k, \quad z_{\vec{\nu}} = z_{\nu_1} \cdots z_{\nu_k}, \\ \{\vec{\nu}\} &= \sum_{j=1}^k \nu'_j, \quad [\vec{\nu}]_- = \sum_{j=1}^k [\nu_j]_-, \quad [\vec{\nu}] = \sum_{j=1}^k [\nu_j]. \end{aligned}$$

Here $\nu'_j = \nu_j - [\nu_j]_-$ for $\nu_j > 0$. From Propositions 3.1 and 3.5 (a) and the resolvent equation $R_0(z) = (1 - \tilde{R}_0(z)W)^{-1}\tilde{R}_0(z)$, we obtain the following

Proposition 3.6. *The following asymptotic expansion holds for z near 0 with $\Im z > 0$.*

(a). *Let $N \in \mathbb{N}$ and $s > 2N + 1$. Then there exists $N_0 \in \mathbb{N}$ depending on N and $\min \sigma_\infty$ such that*

$$R_0(z) = \sum_{j=0}^N z^j R_j + \sum_{\{\vec{\nu}\} + j \leq N}^{(1)} z_{\vec{\nu}} z^j R_{\vec{\nu},j} + R_0^{(N)}(z), \quad (3.32)$$

in $\mathcal{L}(-1, s; 1, -s)$. Here the notation $\sum_{\{\vec{\nu}\} + j \leq N}^{(l)}$ means the finite sum taken over all $\{\vec{\nu}\} \in \sigma_N^k$, $k \geq l$, and $j \geq [\vec{\nu}]_-$ such that $\{\vec{\nu}\} + j \leq N$,

$$R_0 = AF_0; \quad R_1 = AF_1 A^*; \quad (3.33)$$

$$R_{\vec{\nu},0} = AG_{\nu_1, \delta_{\nu_1}} \pi_{\nu_1} W A G_{\nu_2, \delta_{\nu_2}} \pi_{\nu_2} W \cdots A G_{\nu_k, \delta_{\nu_k}} \pi_{\nu_k} A^* \quad (3.34)$$

for $\vec{\nu} = (\nu_1, \nu_2, \dots, \nu_k)$ with $A = (1 - F_0 W)^{-1}$. In particular, if $k = 1$ and $\vec{\nu} = \nu_1$, one has

$$R_{\vec{\nu},0} = A G_{\nu_1, \delta_{\nu_1}} \pi_{\nu_1} A^*. \quad (3.35)$$

R_j (resp. $R_{\vec{\nu},j}$) are in $\mathcal{L}(-1, s; 1, -s)$ for $s > 2j + 1$ (resp. for $s > 2j + \{\vec{\nu}\} + 1$), and $R_0^{(N)}(z) = O(|z|^{N+\epsilon})$ in $\mathcal{L}(-1, s; 1, -s)$, $s > 2N + 1$.

(b). *The above expansion (3.32) can be differentiated in z and one has the estimate*

$$\frac{d}{dz} R_0^{(N)}(z) = O(|z|^{N-1+\epsilon}), \quad (3.36)$$

in $\mathcal{L}(-1, s; 1, -s)$, $s > 2N + 1$, with some $\epsilon > 0$.

For the operator P , zero may be an eigenvalue or a resonance of P . The multiplicity of zero resonance may be large, but does not exceed the sum of multiplicities of the eigenvalues λ of $-\Delta_{\mathbb{S}^{n-1}} + q(\theta)$ such that $\nu = \sqrt{\lambda + \frac{(n-2)^2}{4}} \in]0, 1]$. The existence of an asymptotic expansion of $R(z)$ can be obtained as in [33] by regarding P as perturbation of P_0 . In order to obtain a more “concise” expansion, we regard P as perturbation of \tilde{P}_0 . Let

$$0 < \varsigma_1 < \cdots < \varsigma_{\kappa_0} \leq 1 \quad (3.37)$$

be the points in σ_1 such that P has m_{ς_j} linearly independent ς_j -resonant states with $\sum_{j=1}^{\kappa_0} m_{\varsigma_j} = \mu_r$, μ_r being the multiplicity of zero resonance of P . Modifying the basis of the eigenfunctions $\{\varphi_\nu^{(j)}\}$ of $-\Delta_{\mathbb{S}^{n-1}} + q(\theta)$ if necessary, one can show (see (3.22) and the remarks following it) that there exists a basis of ς_j -resonant states, $u_j^{(i)}$, $i = 1, \dots, m_{\varsigma_j}$ verifying

$$|c_{\varsigma_j}|^{1/2} \langle \widetilde{W} u_j^{(l)}, |x|^{-\frac{n-2}{2} + \varsigma_j} \varphi_{\varsigma_j}^{(l')} \rangle = \delta_{ll'}, \quad 1 \leq l \leq m_{\varsigma_j}, 1 \leq l' \leq m_{\varsigma_j}, \quad 1 \leq j \leq \kappa_0, \quad (3.38)$$

where c_ν is the coefficient of G_{ν, δ_ν} given by

$$c_\nu = -\frac{e^{-i\pi\nu}\Gamma(1-\nu)}{\nu 2^{2\nu+1}\Gamma(1+\nu)}, \nu \in]0, 1[; \quad c_1 = -\frac{1}{8}. \quad (3.39)$$

(see (3.17) and (3.18)) and $\delta_{ll'} = 1$ if $l = l'$; 0 otherwise. As seen in [33], we can choose a basis $\{\phi_j; j = 1, \dots, \mu\}$ of \mathcal{N} , $\mu = \dim \mathcal{N}$, such that

$$\langle \phi_i, \widetilde{W} \phi_j \rangle = \delta_{ij}.$$

Without loss, we can assume that for $1 \leq j \leq \mu_r$, ϕ_j is resonant state of P , while for $\mu_r + 1 \leq j \leq \mu$, ϕ_j is an eigenfunction of P . Define

$$Q_r = \sum_{j=1}^{\mu_r} \langle \widetilde{W} \phi_j, \cdot \rangle \phi_j \quad (3.40)$$

$$Q_e = \sum_{j=\mu_r+1}^{\mu} \langle \widetilde{W} \phi_j, \cdot \rangle \phi_j. \quad (3.41)$$

The following result can be proved by studying an appropriate Grushin problem for $(1 - \widetilde{R}_0(z)\widetilde{W})$ as in [33]. See [14] for the details.

Theorem 3.7. *Let $\mu = \dim \mathcal{N} \neq 0$ and $N \in \mathbb{N}$. Assume $\rho_0 > \max\{4N-2, 2N+4\}$. One has the following asymptotic expansion for $R(z)$ in $\mathcal{L}(-1, s; 1, -s)$, $s > \max\{2N+1, 2\}$:*

$$R(z) = \sum_{j=0}^{N-1} z^j T_j + \sum_{\{\vec{\nu}\}+j \leq N-1}^{(1)} z_{\vec{\nu}} z^j T_{\vec{\nu}, j} + T_e(z) + T_r(z) + T_{er}(z) + O(|z|^{N-1+\epsilon}) \quad (3.42)$$

Here T_j (resp., $T_{\vec{\nu}, j}$) is in $\mathcal{L}(1, -s; -1, s)$ for $s > 2j+1$ (resp., for $s > 2j+1 + \{\vec{\nu}\}$),

$$T_0 = \widetilde{A}F_0, \quad T_1 = \widetilde{A}F_1(1 + \widetilde{W}\widetilde{A}F_0)$$

with $\widetilde{A} = (\Pi'(1 - F_0\widetilde{W})\Pi')^{-1}\Pi'$, Π' is the projection from $H^{1,-s}$ onto $\text{ran}(1 - F_0\widetilde{W})$ corresponding to the decomposition $H^{1,-s} = \ker(1 - F_0\widetilde{W}) \oplus \text{ran}(1 - F_0\widetilde{W})$. The sum $\sum_{\{\vec{\nu}\}+j \leq N}^{(1)}$ has the same meaning as in (3.32) and the first term in this sum is z_{ν_0} with coefficient $T_{\nu_0, 0}$ given by

$$T_{\nu_0, 0} = \widetilde{A}G_{\nu_0, \delta_{\nu_0}}\pi_{\nu_0}(1 + \widetilde{W}\widetilde{A}F_0),$$

where ν_0 is the smallest value of $\nu \in \sigma_\infty$. $T_e(z)$, $T_r(z)$ describe the contributions up to the order $O(|z|^{N-1+\epsilon})$ from eigenfunctions and resonant states, respectively, and $T_{er}(z)$ the interaction between eigenfunctions and resonant states. One has

$$\begin{aligned} T_e(z) &= -z^{-1}\Pi_0 + \sum_{j, \{\vec{\nu}\}+j \leq N-1}^{(-)} z_{\vec{\nu}} z^j T_{e;\vec{\nu};j} \\ T_r(z) &= \sum_{j=1}^{\kappa_0} z_{\zeta_j}^{-1} (\Pi_{r,j} + \sum_{\alpha, \beta, \vec{\nu}, l}^{+,N} z_{\vec{\nu}} z^{|\beta|} (z_{\zeta}^{-\alpha-\beta} z^l T_{r;\vec{\nu}, \alpha, \beta, l, j}), \quad \text{with} \\ \Pi_{r,j} &= e^{i\pi\zeta'_j} \sum_{l=1}^{m_{\zeta_j}} \langle \cdot, u_j^{(l)} \rangle u_j^{(l)}, \quad j = 1, \dots, \kappa_0, \\ T_{er}(z) &= \sum_{j=1}^{\kappa_0} z_{\zeta_j}^{-1} (\Pi_0 \widetilde{W} Q_e F_1 \widetilde{W} \Pi_{r,j} + \Pi_{r,j} \widetilde{W} Q_r F_1 \widetilde{W} \Pi_0 \\ &\quad + \sum_{\alpha, \beta, \vec{\nu}, l}^{+,N} z_{\vec{\nu}} z^{|\beta|} (z_{\zeta}^{-\alpha-\beta} z^l T_{er;\vec{\nu}, \alpha, \beta, l, j}). \end{aligned}$$

Here $\zeta'_j = \zeta_j - [\zeta_j]$, Π_0 is the spectral projection of P at 0, and $T_e(z)$ is of rank not exceeding $\text{Rank } \Pi_0$ with leading singular parts given by $\nu_j \in \sigma_2$:

$$T_{e;\vec{\nu};-1} = (-1)^{k'+1} (\Pi_0 \widetilde{W} G_{\nu_1, 1+\delta_{\nu_1}} \pi_{\nu_1} \widetilde{W}) \cdots (\Pi_0 \widetilde{W} G_{\nu_{k'}, 1+\delta_{\nu_{k'}}} \pi_{\nu_{k'}} \widetilde{W}) \Pi_0, \quad (3.43)$$

for $\vec{\nu} = (\nu_1, \dots, \nu_{k'}) \in \sigma_2^{k'}$ with $\{\vec{\nu}\} \leq 1$, $(z_{\zeta})^{-\alpha} = (z_{\zeta_1})^{-\alpha_1} \cdots (z_{\zeta_{\kappa_0}})^{-\alpha_{\kappa_0}}$. The summation $\sum_{j, \{\vec{\nu}\}+j \leq N-1}^{(-)}$ is taken over all indices $\vec{\nu} \in (\sigma_N)^k$, $k \in \mathbb{N}^*$ and all $j \in \mathbb{Z}$ with $j \geq [\vec{\nu}]_- - 1$ such that $\{\vec{\nu}\} + j \leq N-1$, and the summation $\sum_{\alpha, \beta, \vec{\nu}, l}^{+,N}$ is taken over all possible $\alpha, \beta \in \mathbb{N}^{\kappa_0}$ with $1 \leq |\alpha| \leq N_0$, $|\beta| \geq 1$, $\vec{\nu} = (\nu_1, \dots, \nu_{k'}) \in \sigma_N^{k'}$, $k' \geq 2|\alpha|$, for which there are at least α_k values of ν_j 's belonging to σ_1 with $\nu_j \geq \zeta_k$, for $1 \leq k \leq \kappa_0$, $l \in \mathbb{N}$, satisfying

$$|\beta| + \{\vec{\nu}\} + l - \sum_{k=1}^{\kappa_0} (\alpha_k + \beta_k) \zeta_k \leq N.$$

This theorem is proved [33] for the Schrödinger operator of the form $P = \widetilde{P}_0 + V$ with V bounded and satisfying $|V| \leq C \langle x \rangle^{-\rho_0}$. In the present work, $V = -\frac{\chi_1^2}{r^2} q(\theta) + \widetilde{W}$ has a critical singularity at 0. So we can not directly use the result of [33], but with the help of Proposition 3.5, one can follow the same line of proof to obtain Theorem 3.7 (see [14]). With the sign correction on c_1 and the formula (4.32) in [33], one sees that a 1-resonant state $u_j^{(l)}$ gives rise to a singularity of the leading term $\frac{1}{z \ln z} \langle \cdot, u_j^{(l)} \rangle u_j^{(l)}$ instead of $-\frac{1}{z \ln z} \langle \cdot, u_j^{(l)} \rangle u_j^{(l)}$ as stated in Theorem 4.6 of [33].

4. GENERALIZED RESIDUE OF THE TRACE FUNCTION

Let f be a function satisfying the condition of Theorem 2.1. As we have seen before, $(R(z) - R_0(z))f(P)$ is of trace class for $z \notin \sigma(P)$. Moreover, the application $z \rightarrow T(z) = \text{Tr} [(R(z) - R_0(z))f(P)]$ is meromorphic on $\mathbb{C} \setminus \mathbb{R}_+$. The goal of this section is to calculate the generalized residue of $T(z)$ at $z = 0$, in the sense of Section 2.

First, let us recall some well-known results for the trace ideals $\mathcal{S}_p \subset \mathcal{L}(L^2(\mathbb{R}^n))$, (see [28, IX.4] for details).

Definition 4.1. For $1 \leq p < \infty$, we say that a compact operator $A \in \mathcal{S}_p$ if $|A|^p$ is a trace class operator, where $|A| = \sqrt{A^*A}$, and we set $\|A\|_p = (\text{Tr } |A|^p)^{1/p}$. For $p = \infty$, \mathcal{S}_∞ is the set of the compact operators with $\|A\|_\infty = \|A\|$.

We have the following properties :

Proposition 4.2. ([28, IX.4]) Let $1 \leq p \leq \infty$ and $p^{-1} + q^{-1} = 1$.

- (a) If $A \in \mathcal{S}_p$ and $B \in \mathcal{S}_q$, then $AB \in \mathcal{S}_1$ and $\|AB\|_1 \leq \|A\|_p \cdot \|B\|_q$.
- (b) \mathcal{S}_p is a Banach space with norm $\|\cdot\|_p$.
- (c) $\mathcal{S}_1 \subset \mathcal{S}_p$.
- (d) If $A \in \mathcal{S}_p$, then $A^* \in \mathcal{S}_p$ and $\|A^*\|_p = \|A\|_p$.

By Proposition 3.6, for z near 0 with $\Im z > 0$, one has for $s > 1$:

$$R_0(z) = R_0 + R_0^{(0)}(z) \text{ in } \mathcal{L}(-1, s; 1, -s), \quad (4.1)$$

with $R_0^{(0)}(z) = O(|z|^\epsilon)$. We deduce the following result :

Lemma 4.3. For $m > n/2$, $s > 3$ and z near 0, $\Im z > 0$, $\langle x \rangle^{-s} R_0(z) \langle x \rangle^{-s} \in \mathcal{S}_m$ and there exists a constant C independent of z , such that

$$\|\langle x \rangle^{-s} R_0(z) \langle x \rangle^{-s}\|_m \leq C.$$

Moreover, $\langle x \rangle^{-s} (R(z) - R_0) \langle x \rangle^{-s} = O(|z|^\epsilon)$ in \mathcal{S}_m .

Proof. Let $\chi \in C_0^\infty(\mathbb{R})$ such that $\chi(r) = 1$ for $|r| < 1$. Then $\langle x \rangle^{-s} R_0(z) \langle x \rangle^{-s}$ can be written as

$$\langle x \rangle^{-s} R_0(z) \langle x \rangle^{-s} = F_1(z) + F_2(z),$$

with

$$F_1(z) = \langle x \rangle^{-s} R_0(z) \chi(P_0) \langle x \rangle^{-s}; \quad F_2(z) = \langle x \rangle^{-s} R_0(z) (1 - \chi(P_0)) \langle x \rangle^{-s}.$$

First, let us study $F_2(z)$. Using the resolvent identity, we can decompose $F_2(z) = F_{21} + F_{22}(z)$, where

$$\begin{aligned} F_{21} &= \langle x \rangle^{-s} R_0(-1) (1 - \chi(P_0)) \langle x \rangle^{-s}, \\ F_{22}(z) &= (1 + z) (\langle x \rangle^{-s} R_0(z) \langle x \rangle^{-s'}) (\langle x \rangle^{s'} R_0(-1) (1 - \chi(P_0)) \langle x \rangle^{-s}). \end{aligned}$$

It is easy to check that F_{21} and $\langle x \rangle^{s'} R_0(-1) (1 - \chi(P_0)) \langle x \rangle^{-s}$ are in \mathcal{S}_m , if $s' > 1$ is chosen close to 1. By (4.1), it is clear that $\langle x \rangle^{-s} R_0(z) \langle x \rangle^{-s'}$ is uniformly bounded, for z near 0, $\Im z > 0$. Then we deduce that $F_2(z) \in \mathcal{S}_m$, and $\|F_2(z)\|_m \leq C$ with some constant C independent of z .

Now, let us study $F_1(z)$. We write $F_1(z) = \langle x \rangle^{-s} R_0(z) \langle x \rangle^{-s'} (\langle x \rangle^{s'} \chi(P_0) \langle x \rangle^{-s})$ with $s' > 1$ close to 1. Using a similar argument as above, we can get that $\|F_1(z)\|_m \leq C$ with some constant C independent of z .

Therefore $\langle x \rangle^{-s} R_0(z) \langle x \rangle^{-s} \in \mathcal{S}_m$, and for some C independent of z ,

$$\|\langle x \rangle^{-s} R_0(z) \langle x \rangle^{-s}\|_m \leq \|F_1(z)\|_m + \|F_2(z)\|_m \leq C.$$

Repeating the same arguments with $R_0(z)$ replaced by $R_0(z) - R_0$, we obtain that $\langle x \rangle^{-s} (R(z) - R_0) \langle x \rangle^{-s} = O(|z|^\epsilon)$ in \mathcal{S}_m . \square

We can now establish the main result of this section :

Theorem 4.4. *Assume that $\rho > \max\{6, n+2\}$ and f satisfies the condition of Theorem 2.1. Then the generalized residue of $T(z) = \text{Tr} [(R(z) - R_0(z))f(P)]$ at $z = 0$ is given by*

$$J_0 = \mathcal{N}_0 + \sum_{j=1}^{\kappa_0} \varsigma_j m_j, \quad (4.2)$$

where \mathcal{N}_0 is the multiplicity of zero as the eigenvalue of P and m_j the multiplicity of ς_j -resonance of zero.

Proof. The proof is rather long and is divided in two steps. We write $f \approx g$ if $f(z) - g(z) = O(|z|^{-1+\epsilon})$ for some $\epsilon > 0$, with z in a neighborhood of 0 and $\Im z > 0$.

We fix $k \in \mathbb{N}$ with $k > \frac{n}{2} + 1$ and we use the following resolvent identities :

$$\begin{aligned} R(z) - R_0(z) &= -R_0(z)V R(z) \\ &= \sum_{j=1}^{k-1} (-1)^j (R_0(z)V)^j R_0(z) + (-1)^k R(z)(V R_0(z))^k. \end{aligned}$$

Decompose $T(z) = -\text{Tr} [R_0(z)V f(P)R(z)]$ as $T(z) = T_1(z) + T_2(z)$ where

$$T_1(z) = -\text{Tr} [R_0(z)V f(P)(R_0(z) + \sum_{j=1}^{k-1} (-1)^j (R_0(z)V)^j R_0(z))]. \quad (4.3)$$

$$T_2(z) = (-1)^{k+1} \text{Tr} [R_0(z)V f(P)R(z)(V R_0(z))^k]. \quad (4.4)$$

First, let us study $T_1(z)$. Using the cyclicity of the trace, one has, for $s > 1$, close to 1,

$$\begin{aligned} \text{Tr} [R_0(z)V f(P)R_0(z)] &= \text{Tr} [R_0^2(z)V f(P)] \\ &= \text{Tr} [(\langle x \rangle^{-s} \frac{d}{dz} R_0(z) \langle x \rangle^{-s}) (\langle x \rangle^s V f(P) \langle x \rangle^s)]. \end{aligned}$$

Since $\rho > n + 2$, $\langle x \rangle^s V f(P) \langle x \rangle^s \in \mathcal{S}_1$. So, by Proposition 3.6, we deduce that

$$\text{Tr} [R_0(z)V f(P)R_0(z)] \approx 0. \quad (4.5)$$

Similarly, we can get that the other terms of $T_1(z)$ give the same contribution. Therefore, we have obtained :

$$T_1(z) \approx 0. \quad (4.6)$$

Now, let us study $T_2(z)$. Since $k > \frac{n}{2} + 1 > \frac{n}{2} - 2$, we can define $T_3(z)$ for $z \notin \sigma(P)$ by :

$$T_3(z) = (-1)^{k+1} \text{Tr} [R_0(z)V R(z)(V R_0(z))^k]. \quad (4.7)$$

Then, for $s > 1$ close to 1, we have as previously,

$$\begin{aligned} T_2(z) - T_3(z) &= (-1)^k \text{Tr} [R_0(z)V R(z)(1 - f(P))(V R_0(z))^k] \\ &= (-1)^k \text{Tr} [(\langle x \rangle^{-s} \frac{d}{dz} R_0(z) \langle x \rangle^{-s}) \langle x \rangle^s V (R(z)(1 - f(P))) \\ &\quad (V R_0(z))^{k-1} V \langle x \rangle^s]. \end{aligned}$$

Since $1 - f(t)$ is equal to 0 for t near 0, $R(z)(1 - f(P))$ is uniformly bounded for z in a neighborhood of 0. Moreover, since $\rho_0 > 6$ and $k > \frac{n}{2} + 1$, Proposition 4.2 (a) and Lemma 4.3 imply that, for z near 0 and $z \notin \mathbb{R}_+$,

$$\| (VR_0(z))^{k-1} V \langle x \rangle^s \|_1 = O(1). \quad (4.8)$$

By Proposition 3.6 (b), we conclude that, for some $\epsilon > 0$ and z near 0 with $\Im z > 0$,

$$T_2(z) - T_3(z) \approx 0. \quad (4.9)$$

The main task of the proof is to estimate $T_3(z)$ at 0.

• **Step 1:** Assume that 0 is not an eigenvalue of P .

Using the cyclicity of the trace, we remark that

$$T_3(z) = (-1)^{k+1} \text{Tr} \left[\left(\frac{d}{dz} R_0(z) \right) V R(z) V (R_0(z) V)^{k-1} \right]. \quad (4.10)$$

Let us introduce the following notation :

$$U_1 = \text{sgn}(V) |V|^{\frac{1}{2}}, \quad U_2 = |V|^{\frac{1}{2}},$$

$$S_1(z) = U_1 \frac{d}{dz} R_0(z) U_2, \quad S_2(z) = U_1 R(z) U_2, \quad S_3(z) = U_1 R_0(z) U_2.$$

Then,

$$T_3(z) = (-1)^{k+1} \text{Tr} [S_1(z) S_2(z) S_3^{k-1}(z)]. \quad (4.11)$$

As previously, by Lemma 4.3, we have for $k > \frac{n}{2} + 1$, $\rho_0 > 6$ and z near 0 with $\Im(z) > 0$,

$$\|S_3^{k-1}(z)\|_1 = O(1). \quad (4.12)$$

By Proposition 3.6 (b), with $N = 1$,

$$\frac{d}{dz} R_0(z) = R_1 + \sum_{\{\vec{\nu}\} \leq 1}^{(1)} \frac{d}{dz} z_{\vec{\nu}} R_{\vec{\nu},0} + O(|z|^\epsilon), \quad (4.13)$$

in $\mathcal{L}(-1, s; 1, -s)$, for $s > 3$ and some $\epsilon > 0$. Since $\rho_0 > 6$, we deduce that

$$S_1(z) = S_{11}(z) + S_{12}(z), \quad (4.14)$$

with

$$S_{11}(z) = U_1 R_1 U_2 + \sum_{\{\vec{\nu}\} \leq 1}^{(1)} \frac{d}{dz} z_{\vec{\nu}} U_1 R_{\vec{\nu},0} U_2 \quad (4.15)$$

and

$$S_{12}(z) = O(|z|^\epsilon) \quad \text{in } \mathcal{L}(L^2(\mathbb{R}^n)). \quad (4.16)$$

In the same way, we can use Theorem 3.7, with $N = 1$ (and $\rho_0 > 6$) to decompose $S_2(z)$. Note that $T_e(z) = 0$, $T_{er}(z) = 0$ since 0 is not an eigenvalue of P . So we have,

$$S_2(z) = S_{21}(z) + S_{22}(z) + S_{23}(z) \quad (4.17)$$

with

$$S_{21}(z) = \sum_{j=1}^{\kappa_0} z_{\zeta_j}^{-1} U_1 \Pi_{r,j} U_2, \quad (4.18)$$

$$S_{22}(z) = \sum_{j=1}^{\kappa_0} \sum_{\alpha, \beta, \vec{\nu}, l}^{+,1} z_{\zeta_j}^{-1} z_{\vec{\nu}} z^{|\beta|} (z_{\zeta_j})^{-\alpha-\beta} z^l U_1 T_{r;\vec{\nu}, \alpha, \beta, l, j} U_2, \quad (4.19)$$

$$S_{23}(z) = O(1) \quad \text{in } \mathcal{L}(L^2(\mathbb{R}^n)). \quad (4.20)$$

It follows from the previous discussion that

$$T_3(z) = T_{31}(z) + T_{32}(z) + T_{33}(z) \quad (4.21)$$

where

$$T_{31}(z) = (-1)^{k+1} \operatorname{Tr} [S_{11}(z) S_{21}(z) S_3^{k-1}(z)], \quad (4.22)$$

$$T_{32}(z) = (-1)^{k+1} \operatorname{Tr} [S_{11}(z) S_{22}(z) S_3^{k-1}(z)], \quad (4.23)$$

$$T_{33}(z) = (-1)^{k+1} \operatorname{Tr} [(S_{11}(z) S_{23}(z) + S_{12}(z) S_{21}(z) + S_{12}(z) S_{22}(z) + S_{12}(z) S_{23}(z)) S_3^{k-1}(z)]. \quad (4.24)$$

First, let us study $T_{33}(z)$. We have the following result :

Lemma 4.5.

$$T_{33}(z) \approx 0. \quad (4.25)$$

Proof. Clearly, we have for some $\epsilon > 0$,

$$S_{11}(z) = O(|z|^{-1+\epsilon}) \quad , \quad S_{21}(z) = O(|z \ln z|^{-1}) \quad \text{in } \mathcal{L}(L^2(\mathbb{R}^n)). \quad (4.26)$$

In the same way, we want to estimate $S_{22}(z)$. We note that the summation $\sum_{\alpha, \beta, \vec{\nu}, l}^{+,1}$ is taken over all possible $\alpha, \beta \in \mathbb{N}^{\kappa_0}$ with $1 \leq |\alpha| \leq N_0$, $|\beta| \geq 1$, $\vec{\nu} = (\nu_1, \dots, \nu_{k'}) \in \sigma_2^{k'}$, $k' \geq 2|\alpha|$, for which there are at least α_k values of ν_j 's belonging to σ_1 with $\nu_j \geq \varsigma_k$, for $1 \leq k \leq \kappa_0$ and $l \in \mathbb{N}$, satisfying the condition

$$|\beta| + \{\vec{\nu}\} + l - \sum_{k=1}^{\kappa_0} (\alpha_k + \beta_k) \varsigma_k \leq 1.$$

It follows that

$$z^{|\beta|} (z_{\zeta_j})^{-\beta} = O(1) \quad , \quad z_{\vec{\nu}} (z_{\zeta_j})^{-\alpha} = O(|z|^\epsilon). \quad (4.27)$$

We emphasize that in the second estimate of (4.27), we have used $k' \geq 2|\alpha|$. Then, we have obtained,

$$S_{22}(z) \approx 0 \quad \text{in } \mathcal{L}(L^2(\mathbb{R}^n)). \quad (4.28)$$

Then, the lemma follows from (4.12), (4.16), (4.20), (4.26) and (4.28). \square

The following elementary lemma will be useful to estimate $T_{31}(z)$ and $T_{32}(z)$.

Lemma 4.6. *Let $u \in \mathcal{N}(P)$, $A = (1 - F_0 W)^{-1}$. Then*

$$A^* V u = -\widetilde{W} u \quad , \quad R_0 V u = -u.$$

Proof. By definition, if $u \in \mathcal{N}(P)$, $F_0 \widetilde{W} u = u$. Thus, $V u = (W - \widetilde{W}) u = -(1 - W F_0) \widetilde{W} u$ which proves the first equality. Using the same argument, if $P u = (P_0 + V) u = 0$ then Proposition 3.6 implies that $R_0 V u = -u$. \square

Now, let us study $T_{31}(z)$; we shall see that $T_{31}(z)$ gives the leading term of the generalized residue. We denote δ_{ij} the Kronecker symbol : $\delta_{ij} = 1$ if $i = j$, $\delta_{ij} = 0$ otherwise.

Lemma 4.7.

$$T_{31}(z) \approx -\frac{1}{z} \sum_{j=1}^{\kappa_0} \varsigma_j m_{\varsigma_j} - \frac{\delta_{1, \varsigma_{\kappa_0}}}{z \ln z} \left(m_1 - \sum_{l=1}^{m_1} \langle R_1 V u_{\kappa_0}^{(l)}, V u_{\kappa_0}^{(l)} \rangle \right). \quad (4.29)$$

Proof. It follows from (4.22) that

$$\begin{aligned} (-1)^{k+1} T_{31}(z) &= \text{Tr} \left[U_1 R_1 V \sum_{j=1}^{\kappa_0} z_{\varsigma_j}^{-1} \Pi_{r,j} U_2 S_3^{k-1}(z) \right] \\ &+ \text{Tr} \left[\sum_{\{\vec{\nu}\} \leq 1}^{(1)} \frac{d}{dz} z_{\vec{\nu}} U_1 R_{\vec{\nu},0} V \sum_{j=1}^{\kappa_0} z_{\varsigma_j}^{-1} \Pi_{r,j} U_2 S_3^{k-1}(z) \right] \\ &= (a) + (b). \end{aligned}$$

First, let us study (b). As we will see, the cases $\varsigma_{\kappa_0} < 1$ and $\varsigma_{\kappa_0} = 1$ are slightly different.

◦ *Case 1 :* Assume that $\varsigma_{\kappa_0} < 1$.

It follows that $\varsigma_j' = \varsigma_j$, and for $j = 1, \dots, \kappa_0$,

$$\Pi_{r,j} = e^{i\pi\varsigma_j} \sum_{l=1}^{m_{\varsigma_j}} \langle \cdot, u_j^{(l)} \rangle u_j^{(l)}.$$

We deduce that

$$\begin{aligned} (b) &= \sum_{j=1}^{\kappa_0} \sum_{l=1}^{m_{\varsigma_j}} \sum_{\{\vec{\nu}\} \leq 1}^{(1)} e^{i\pi\varsigma_j} z_{\varsigma_j}^{-1} \frac{d}{dz} z_{\vec{\nu}} \text{Tr} [U_1 R_{\vec{\nu},0} V \langle U_2 S_3^{k-1}(z) \cdot, u_j^{(l)} \rangle u_j^{(l)}] \\ &= \sum_{j=1}^{\kappa_0} \sum_{l=1}^{m_{\varsigma_j}} \sum_{\{\vec{\nu}\} \leq 1}^{(1)} e^{i\pi\varsigma_j} z_{\varsigma_j}^{-1} \frac{d}{dz} z_{\vec{\nu}} \text{Tr} [U_1 R_{\vec{\nu},0} V \langle \cdot, (S_3^{k-1}(z))^* U_2 u_j^{(l)} \rangle u_j^{(l)}] \\ &= \sum_{j=1}^{\kappa_0} \sum_{l=1}^{m_{\varsigma_j}} \sum_{\{\vec{\nu}\} \leq 1}^{(1)} e^{i\pi\varsigma_j} z_{\varsigma_j}^{-1} \frac{d}{dz} z_{\vec{\nu}} \langle U_1 R_{\vec{\nu},0} V u_j^{(l)}, (S_3^{k-1}(z))^* U_2 u_j^{(l)} \rangle \\ &= \sum_{j=1}^{\kappa_0} \sum_{l=1}^{m_{\varsigma_j}} \sum_{\{\vec{\nu}\} \leq 1}^{(1)} e^{i\pi\varsigma_j} z_{\varsigma_j}^{-1} \frac{d}{dz} z_{\vec{\nu}} \langle R_{\vec{\nu},0} V u_j^{(l)}, U_1 (S_3^{k-1}(z))^* U_2 u_j^{(l)} \rangle. \end{aligned}$$

Using (3.34) with $\vec{\nu} = (\nu_1, \dots, \nu_p)$, we have

$$R_{\vec{\nu},0} V u_j^{(l)} = A G_{\nu_1, \delta_{\nu_1}} \pi_{\nu_1} W A G_{\nu_2, \delta_{\nu_2}} \pi_{\nu_2} W \cdots A G_{\nu_p, \delta_{\nu_p}} \pi_{\nu_p} A^* V u_j^{(l)}. \quad (4.30)$$

Thus, Lemma 4.6 implies

$$R_{\vec{\nu},0} V u_j^{(l)} = -A G_{\nu_1, \delta_{\nu_1}} \pi_{\nu_1} W A G_{\nu_2, \delta_{\nu_2}} \pi_{\nu_2} W \cdots A G_{\nu_p, \delta_{\nu_p}} \pi_{\nu_p} \widetilde{W} u_j^{(l)}. \quad (4.31)$$

It is easy to see that for $\vec{\nu} = (\nu_1, \dots, \nu_p)$ with $\sum_{i=1}^p \nu_i > \varsigma_j$, one has $z_{\varsigma_j}^{-1} \frac{d}{dz} z_{\vec{\nu}} = O(|z|^{-1+\epsilon})$ for some $\epsilon > 0$. It follows from (4.12) that it suffices to evaluate (b) for $\vec{\nu} = (\nu_1, \dots, \nu_p)$ with $\sum_{i=1}^p \nu_i \leq \varsigma_j$,

◇ Assume first $\exists i \in \{1, \dots, p\}$ such that $\nu_i = \varsigma_j$.

One deduces that $p = 1$, i.e $\vec{\nu} = \nu_1$. Moreover, since $\varsigma_{\kappa_0} < 1$, one has $\nu_1 = \varsigma_j < 1$. Using (3.35), we obtain

$$R_{\vec{\nu},0} V u_j^{(l)} = -A G_{\varsigma_j,0} \pi_{\varsigma_j} \widetilde{W} u_j^{(l)}. \quad (4.32)$$

Recalling that

$$\pi_{\varsigma_j} = \sum_{l'=1}^{n_{\varsigma_j}} (\cdot, \varphi_{\varsigma_j}^{(l')}) \otimes \varphi_{\varsigma_j}^{(l')},$$

and using (3.17) and (3.39), we deduce

$$\begin{aligned} G_{\varsigma_j,0} \pi_{\varsigma_j} \widetilde{W} u_j^{(l)} &= \sum_{l'=1}^{n_{\varsigma_j}} \int_0^{+\infty} G_{\varsigma_j,0}(r, \tau) (\widetilde{W} u_j^{(l)}, \varphi_{\varsigma_j}^{(l')}) \varphi_{\varsigma_j}^{(l')} \tau^{n-1} d\tau \\ &= c_{\varsigma_j} r^{-\frac{n-2}{2}+\varsigma_j} \sum_{l'=1}^{n_{\varsigma_j}} \langle \widetilde{W} u_j^{(l)}, |x|^{-\frac{n-2}{2}+\varsigma_j} \varphi_{\varsigma_j}^{(l')} \rangle \varphi_{\varsigma_j}^{(l')}. \end{aligned} \quad (4.33)$$

It follows from the normalization condition of resonant states given in (3.38) that

$$G_{\varsigma_j,0} \pi_{\varsigma_j} \widetilde{W} u_j^{(l)} = -e^{-i\pi\varsigma_j} |c_{\varsigma_j}|^{\frac{1}{2}} r^{-\frac{n-2}{2}+\varsigma_j} \varphi_{\varsigma_j}^{(l)}. \quad (4.34)$$

Thus, in this case, we have

$$R_{\vec{\nu},0} V u_j^{(l)} = e^{-i\pi\varsigma_j} |c_{\varsigma_j}|^{\frac{1}{2}} A \left(|x|^{-\frac{n-2}{2}+\varsigma_j} \varphi_{\varsigma_j}^{(l)} \right). \quad (4.35)$$

◇ Assume now that $\forall i \in \{1, \dots, p\}$, $\nu_i < \varsigma_j$.

In particular $\nu_p < 1$ and using (3.24), we obtain

$$G_{\nu_p,0} \pi_{\nu_p} \widetilde{W} u_j^{(l)} = G_{\nu_p,0} \sum_{l'=1}^{n_{\nu_p}} (\widetilde{W} u_j^{(l)}, \varphi_{\nu_p}^{(l')}) \otimes \varphi_{\nu_p}^{(l')} \quad (4.36)$$

$$= c_{\nu_p} r^{-\frac{n-2}{2}+\nu_p} \sum_{l'=1}^{n_{\nu_p}} \langle \widetilde{W} u_j^{(l)}, |x|^{-\frac{n-2}{2}+\nu_p} \varphi_{\nu_p}^{(l')} \rangle \varphi_{\nu_p}^{(l')} \quad (4.37)$$

$$= 0. \quad (4.38)$$

Thus, in this case, using (4.31), we have

$$R_{\vec{\nu},0} V u_j^{(l)} = 0. \quad (4.39)$$

As a conclusion of (4.35) and (4.39), we obtain

$$(b) \approx \sum_{j=1}^{\kappa_0} \sum_{l=1}^{m_{\varsigma_j}} |c_{\varsigma_j}|^{\frac{1}{2}} z_{\varsigma_j}^{-1} \frac{d}{dz} (z_{\varsigma_j}) \langle A(|x|^{-\frac{n-2}{2}+\varsigma_j} \varphi_{\varsigma_j}^{(l)}), U_1(S_3^{k-1}(z))^* U_2 u_j^{(l)} \rangle. \quad (4.40)$$

By Proposition 3.6, $(S_3^{k-1}(z))^* = (U_2 R_0 U_1)^{k-1} + O(|z|^\epsilon)$ in $L^2(\mathbb{R}^n)$. Moreover, since $\varsigma_j < 1$, $z_{\varsigma_j}^{-1} \frac{d}{dz} (z_{\varsigma_j}) = \frac{\varsigma_j}{z}$. Thus,

$$\begin{aligned}
(b) &\approx \frac{1}{z} \sum_{j=1}^{\kappa_0} \sum_{l=1}^{m_{\varsigma_j}} |c_{\varsigma_j}|^{\frac{1}{2}} \varsigma_j \langle A(|x|^{-\frac{n-2}{2}+\varsigma_j} \varphi_{\varsigma_j}^{(l)}), U_1(U_2 R_0 U_1)^{k-1} U_2 u_j^{(l)} \rangle \\
&\approx \frac{1}{z} \sum_{j=1}^{\kappa_0} \sum_{l=1}^{m_{\varsigma_j}} |c_{\varsigma_j}|^{\frac{1}{2}} \varsigma_j \langle A(|x|^{-\frac{n-2}{2}+\varsigma_j} \varphi_{\varsigma_j}^{(l)}), V(R_0 V)^{k-1} u_j^{(l)} \rangle \\
&\approx \frac{(-1)^{k-1}}{z} \sum_{j=1}^{\kappa_0} \sum_{l=1}^{m_{\varsigma_j}} |c_{\varsigma_j}|^{\frac{1}{2}} \varsigma_j \langle A(|x|^{-\frac{n-2}{2}+\varsigma_j} \varphi_{\varsigma_j}^{(l)}), V u_j^{(l)} \rangle \\
&\approx \frac{(-1)^k}{z} \sum_{j=1}^{\kappa_0} \sum_{l=1}^{m_{\varsigma_j}} |c_{\varsigma_j}|^{\frac{1}{2}} \varsigma_j \langle |x|^{-\frac{n-2}{2}+\varsigma_j} \varphi_{\varsigma_j}^{(l)}, \widetilde{W} u_j^{(l)} \rangle,
\end{aligned}$$

where we have used Lemma 4.6 in the two last equations. Using again the normalization condition of resonant states given in (3.38), we obtain :

$$(b) \approx \frac{(-1)^k}{z} \sum_{j=1}^{\kappa_0} \sum_{l=1}^{m_{\varsigma_j}} \varsigma_j = \frac{(-1)^k}{z} \sum_{j=1}^{\kappa_0} m_{\varsigma_j} \varsigma_j. \quad (4.41)$$

◦ *Case 2* : assume that $\varsigma_{\kappa_0} = 1$.

In this case, $\varsigma'_j = \varsigma_j$ for $j = 1, \dots, \kappa_0 - 1$ and $\varsigma'_{\kappa_0} = 0$. So we can write,

$$(b) = \sum_{j=1}^{\kappa_0-1} \sum_{l=1}^{m_{\varsigma_j}} \sum_{\{\vec{\nu}\} \leq 1}^{(1)} e^{i\pi\varsigma_j} z_{\varsigma_j}^{-1} \frac{d}{dz} z_{\vec{\nu}} \langle R_{\vec{\nu},0} V u_j^{(l)}, U_1(S_3^{k-1}(z))^* U_2 u_j^{(l)} \rangle + (b'), \quad (4.42)$$

where we set

$$(b') = \sum_{l=1}^{m_1} \sum_{\{\vec{\nu}\} \leq 1}^{(1)} \frac{1}{z \ln z} \frac{d}{dz} z_{\vec{\nu}} \langle R_{\vec{\nu},0} V u_{\kappa_0}^{(l)}, U_1(S_3^{k-1}(z))^* U_2 u_{\kappa_0}^{(l)} \rangle. \quad (4.43)$$

The first part of (b) in (4.42) can be calculated exactly as in the case 1. Thus, one has

$$(b) \approx \frac{(-1)^k}{z} \sum_{j=1}^{\kappa_0-1} m_{\varsigma_j} \varsigma_j + (b'). \quad (4.44)$$

Now, let us study (b'). Using the same approach, an easy calculus gives

$$(b') \approx (-1)^{k-1} \sum_{l=1}^{m_1} \frac{1}{z \ln z} \frac{d}{dz} (z \ln z) \langle R_{1,0} V u_{\kappa_0}^{(l)}, V u_{\kappa_0}^{(l)} \rangle. \quad (4.45)$$

Using (3.35), we have

$$R_{1,0} V u_{\kappa_0}^{(l)} = A G_{1,1} \pi_1 A^* V u_{\kappa_0}^{(l)} = -A G_{1,1} \pi_1 \widetilde{W} u_{\kappa_0}^{(l)}. \quad (4.46)$$

Recalling that

$$\pi_1 = \sum_{l'=1}^{n_1} (\cdot, \varphi_1^{(l')}) \otimes \varphi_1^{(l')},$$

we deduce

$$\begin{aligned} G_{1,1}\pi_1\widetilde{W}u_{\kappa_0}^{(l)} &= \sum_{l'=1}^{n_1} \int_0^{+\infty} G_{1,1}(r, \tau) (\widetilde{W}u_{\kappa_0}^{(l)}, \varphi_1^{(l')}) \varphi_1^{(l')} \tau^{n-1} d\tau \\ &= c_1 r^{-\frac{n-2}{2}+1} \sum_{l'=1}^{n_1} \langle \widetilde{W}u_{\kappa_0}^{(l)}, |x|^{-\frac{n-2}{2}+1} \varphi_1^{(l')} \rangle \varphi_1^{(l')}, \end{aligned}$$

where we have used (3.18) and (3.39) in the last equation. As previously, it follows from the normalization condition of resonant states that

$$G_{1,1}\pi_1\widetilde{W}u_{\kappa_0}^{(l)} = -|c_1|^{\frac{1}{2}} r^{-\frac{n-2}{2}+1} \varphi_1^{(l)}. \quad (4.47)$$

Thus,

$$R_{1,0}Vu_{\kappa_0}^{(l)} = |c_1|^{\frac{1}{2}} A\left(|x|^{-\frac{n-2}{2}+1} \varphi_1^{(l)}\right). \quad (4.48)$$

It follows that

$$\begin{aligned} (b') &\approx (-1)^{k-1} \sum_{l=1}^{m_1} \frac{1}{z \ln z} \frac{d}{dz} (z \ln z) |c_1|^{\frac{1}{2}} \langle A\left(|x|^{-\frac{n-2}{2}+1} \varphi_1^{(l)}\right), Vu_{\kappa_0}^{(l)} \rangle \\ &\approx (-1)^k \sum_{l=1}^{m_1} \frac{1}{z \ln z} \frac{d}{dz} (z \ln z) |c_1|^{\frac{1}{2}} \langle |x|^{-\frac{n-2}{2}+1} \varphi_1^{(l)}, \widetilde{W}u_{\kappa_0}^{(l)} \rangle \\ &\approx (-1)^k m_1 \left(\frac{1}{z} + \frac{1}{z \ln z} \right), \end{aligned}$$

where we have used again the normalization condition of resonant states. Thus, in this case, we have obtained

$$(b) \approx \frac{(-1)^k}{z} \sum_{j=1}^{\kappa_0-1} m_{\varsigma_j} \varsigma_j + (-1)^k m_1 \left(\frac{1}{z} + \frac{1}{z \ln z} \right). \quad (4.49)$$

As a conclusion, we have proved in all cases

$$(b) \approx \frac{(-1)^k}{z} \sum_{j=1}^{\kappa_0} m_{\varsigma_j} \varsigma_j + \delta_{\varsigma_{\kappa_0}, 1} (-1)^k \frac{m_1}{z \ln z}. \quad (4.50)$$

It remains to study (a). Using the same strategy as for (b), we obtain easily

$$(a) \approx (-1)^{k-1} \sum_{j=1}^{k_0} \sum_{l=1}^{m_{\varsigma_j}} e^{i\pi\varsigma'_j} z_{\varsigma_j}^{-1} \langle R_1 Vu_j^{(l)}, Vu_j^{(l)} \rangle \quad (4.51)$$

If $\varsigma_{\kappa_0} < 1$, it is clear that (a) is negligible. If $\varsigma_{\kappa_0} = 1$, we obtain

$$(a) \approx \frac{(-1)^{k-1}}{z \ln z} \sum_{l=1}^{m_1} \langle R_1 Vu_{\kappa_0}^{(l)}, Vu_{\kappa_0}^{(l)} \rangle, \quad (4.52)$$

and the lemma is proved. \square

For $T_{32}(z)$, we have the following

Lemma 4.8.

$$T_{32}(z) \approx 0. \quad (4.53)$$

Proof. Using (4.26) and (4.28), we see that it suffices to prove that the following term is negligible :

$$(c) = \sum_{j=1}^{\kappa_0} \sum_{\{\vec{\nu}\} \leq 1}^{(1)} \sum_{\alpha, \beta, \vec{\nu}_1, l}^{+,1} z_{\varsigma_j}^{-1} z_{\vec{\nu}_1} z^{|\beta|} (z_{\varsigma_j})^{-\alpha-\beta} z^l \frac{d}{dz} z_{\vec{\nu}} \text{Tr} [U_1 R_{\vec{\nu},0} V T_{r;\vec{\nu}_1, \alpha, \beta, l, j} U_2 S_3^{k-1}(z)].$$

First, let us remark that if $\sum_{i=1}^p \nu_i \geq \varsigma_j$, the same argument as (4.27) implies that

$$z_{\varsigma_j}^{-1} z_{\vec{\nu}_1} z^{|\beta|} (z_{\varsigma_j})^{-\alpha-\beta} z^l \frac{d}{dz} z_{\vec{\nu}} \approx 0. \quad (4.54)$$

So, it remains to estimate (c) when $\vec{\nu} = (\nu_1, \dots, \nu_p)$ satisfies $\sum_{i=1}^p \nu_i < \varsigma_j$. From the proof of Theorem 3.7, we know that $T_{r;\vec{\nu}_1, \alpha, \beta, l, j}$ is a linear combination of

$$\Pi_{r,j} B_{r;\vec{\nu}_1, \alpha, \beta, l, j} \quad \text{and} \quad A_{r;\vec{\nu}_1, \alpha, \beta, l, j} \widetilde{W} G_{\mu, \delta_\mu} \pi_\mu B_{r;\vec{\nu}_1, \alpha, \beta, l, j}, \quad (4.55)$$

where $A_{r;\vec{\nu}_1, \alpha, \beta, l, j}$ is a bounded operator in $\mathcal{L}(1, -s; 1, -s)$, $s > 3$ and $B_{r;\vec{\nu}_1, \alpha, \beta, l, j}$ is a bounded operator in $\mathcal{L}(-1, s; -1, s)$ for $s > 3$.

◦ Let us study the contribution coming from $\Pi_{r,j} B_{r;\vec{\nu}_1, \alpha, \beta, l, j}$.

To simplify the notation, we write $B = B_{r;\vec{\nu}_1, \alpha, \beta, l, j}$. In this case, as previously, we have

$$\text{Tr} [U_1 R_{\vec{\nu},0} V \Pi_{r,j} B U_2 S_3^{k-1}(z)] = \sum_{l=1}^{m_{\varsigma_j}} e^{i\pi \varsigma_j'} \langle R_{\vec{\nu},0} V u_j^{(l)}, U_1 (S_3^{k-1}(z))^* U_2 B^* u_j^{(l)} \rangle. \quad (4.56)$$

Since $\nu_p < \varsigma_j$, as in (4.39), we obtain $R_{\vec{\nu},0} V u_j^{(l)} = 0$.

◦ Now, let us study the contribution coming from $A_{r;\vec{\nu}_1, \alpha, \beta, l, j} \Pi_{r,j} \widetilde{W} G_{\mu, \delta_\mu} \pi_\mu B_{r;\vec{\nu}_1, \alpha, \beta, l, j}$.

We remark that, if $\mu < \varsigma_j$, the same argument as before and (3.24) imply $\Pi_{r,j} \widetilde{W} G_{\mu, \delta_\mu} \pi_\mu = 0$. Thus, it suffices to study the case $\mu \geq \varsigma_j$. To do this, we refer to the proof of Theorem 4.6 in [33]. The term $G_{\mu, \delta_\mu} \pi_\mu B_{r;\vec{\nu}_1, \alpha, \beta, l, j}$ comes from the expansion of $L_1(z) \widetilde{W} (D_0 + D_1(z)) \widetilde{R}_0(z)$ with $D_0 \in \mathcal{L}(1, -s; 1, -s)$ for any $s > 1$ and $D_0|_{\mathcal{N}} = 0$ and

$$\begin{aligned} L_1(z) &= \sum_{\mu \in \sigma_1} z_\mu G_{\mu, \delta_\mu} \pi_\mu \\ D_1(z) &= z D_1 + \sum_{\mu \in \sigma_1} z_\mu D_{\mu,0} \end{aligned}$$

for some D_1 and $D_{\mu,0}$ in $\mathcal{L}(1, -s; 1, -s)$ with $s > 3$, so the coefficient in front of $G_{\mu, \delta_\mu} \pi_\mu B_{r;\vec{\nu}_1, \alpha, \beta, l, j}$ is $O(z^\mu)$. In the same way, the term $A_{r;\vec{\nu}_1, \alpha, \beta, l, j} \Pi_{r,j} \widetilde{W}$ comes from the expansion of

$$(1 - (D_0 + D_1(z)) L_1(z) \widetilde{W}) I_r(z) Q_r - \Pi_r(z) \widetilde{W} Q_r,$$

where $I_r(z)$ is of the form

$$I_r(z) = (1 + O(|z|^\epsilon)) \Pi_r(z) \widetilde{W}, \quad \Pi_r(z) = \sum_{j=1}^{\kappa_0} \frac{1}{z_{\varsigma_j}} \Pi_{r,j}$$

with $\Pi_{r,j}$ defined in Theorem 3.7 (see p. 1931 of [33]), so the coefficient in front of $A_{r;\vec{\nu}_1,\alpha,\beta,l,j}\Pi_{r,j}\widetilde{W}$ is $O(|z|^{-\varsigma_j+\epsilon})$. It follows that the coefficient of $T_{r;\vec{\nu}_1,\alpha,\beta,l,j}$ in Theorem 3.7 is $O(|z|^{\epsilon+\mu-\varsigma_j}) = O(|z|^\epsilon)$. As a conclusion, $(c) = O(|z|^{-1+\epsilon})$ and the lemma is proved. \square

• **Step 2 :** Assume that 0 is an eigenvalue of P .

We follow the same strategy. $S_1(z), S_2(z), S_3(z), S_{11}(z), S_{12}(z)$ are the same as before. Using Theorem 3.7 with $N = 1$, $S_2(z)$ can be decomposed as

$$S_2(z) = U_1 R(z) U_2 = \sum_{j=1}^5 S_{2j}(z), \quad (4.57)$$

with

$$\begin{aligned} S_{21}(z) &= U_1 T_r(z) U_2 \\ S_{22}(z) &= -\frac{1}{z} U_1 \Pi_0 U_2 \\ S_{23}(z) &= \sum_{j, \{\vec{\nu}\}+j \leq 0}^{(-)} z_{\vec{\nu}} z^j U_1 T_{e;\vec{\nu},j} U_2 \\ S_{24}(z) &= U_1 T_{er}(z) U_2 \\ S_{25}(z) &= O(1). \end{aligned}$$

Thus, it follows that

$$T_3(z) = (-1)^{k+1} \text{Tr} [S_1(z) S_2(z) S_3^{k-1}(z)] = \sum_{j=1}^5 T_{3j}(z), \quad (4.58)$$

where for $j = 1, \dots, 4$,

$$T_{3j}(z) = (-1)^{k+1} \text{Tr} [S_{11}(z) S_{2j}(z) S_3^{k-1}(z)], \quad (4.59)$$

and

$$T_{35}(z) = (-1)^{k+1} \text{Tr} [(S_{11}(z) S_{25}(z) + S_{12}(z) S_2(z)) S_3^{k-1}(z)]. \quad (4.60)$$

It follows from (4.12), (4.16) and (4.26) that $T_{35}(z) \approx 0$. Now, let us establish the following lemma which will be useful to estimate the other terms.

Lemma 4.9. *For $\vec{\nu} = (\nu_1, \dots, \nu_p) \in (\sigma_1)^p$, one has :*

$$R_{\vec{\nu},0} V \Pi_0 = 0 \quad , \quad \Pi_0 V R_{\vec{\nu},0} = 0. \quad (4.61)$$

Proof. We only prove the first assertion since the other one is similar. Let Ψ_j , $j = 1, \dots, \mathcal{N}_0$, be an orthonormal basis of the eigenspace of P with eigenvalue 0. As previously, one has

$$R_{\vec{\nu},0} V \Phi_j = -A G_{\nu_1, \delta_{\nu_1}} \pi_{\nu_1} W A G_{\nu_2, \delta_{\nu_2}} \pi_{\nu_2} W \cdots A G_{\nu_p, \delta_{\nu_p}} \pi_{\nu_p} \widetilde{W} \Phi_j, \quad (4.62)$$

and

$$\begin{aligned}
G_{\nu_p, \delta_{\nu_p}} \pi_{\nu_p} \widetilde{W} \Phi_j &= G_{\nu_p, \delta_{\nu_p}} \sum_{l=1}^{n_{\nu_p}} (\widetilde{W} \Phi_j, \varphi_{\nu_p}^{(l)}) \otimes \varphi_{\nu_p}^{(l)}, \\
&= c_{\nu_p} \langle \widetilde{W} \Phi_j, |x|^{-\frac{n-2}{2} + \nu_p} \varphi_{\nu_p}^{(l)} \rangle r^{-\frac{n-2}{2} + \nu_p} \otimes \varphi_{\nu_p}^{(l)}, \\
&= 0,
\end{aligned}$$

where we have used (3.22) in the last equation, with $u = \Phi_j \in L^2(\mathbb{R}^n)$. \square

First, let us study $T_{32}(z)$. We have the following result :

Lemma 4.10.

$$T_{32}(z) \approx -\frac{\mathcal{N}_0}{z}. \quad (4.63)$$

Proof. We can decompose $T_{32}(z)$ as

$$T_{32}(z) = I_1(z) + I_2(z) \quad (4.64)$$

with

$$\begin{aligned}
I_1(z) &= \frac{(-1)^k}{z} \sum_{\{\vec{\nu}\} \leq 1}^{(1)} \frac{d}{dz} z_{\vec{\nu}} \text{Tr} [U_1 R_{\vec{\nu}, 0} V \Pi_0 U_2 S_3^{k-1}(z)], \\
I_2(z) &= \frac{(-1)^k}{z} \text{Tr} [U_1 R_1 V \Pi_0 U_2 S_3^{k-1}(z)].
\end{aligned}$$

By Lemma 4.9, $I_1(z) = 0$. Now, let us study $I_2(z)$. As previously, we have

$$\begin{aligned}
\text{Tr} [U_1 R_1 V \Pi_0 U_2 S_3^{k-1}(z)] &= \sum_{j=1}^{\mathcal{N}_0} \langle R_1 V \Phi_j, U_1 (S_3^{k-1}(z))^* U_2 \Phi_j \rangle \\
&= (-1)^{k-1} \sum_{j=1}^{\mathcal{N}_0} \langle R_1 V \Phi_j, V \Phi_j \rangle + O(|z|^\epsilon).
\end{aligned}$$

Recall that $R_1 = A F_1 A^*$ and $A^* V \Phi_j = -\widetilde{W} \Phi_j$. Thus,

$$\text{Tr} [U_1 R_1 V \Pi_0 U_2 S_3^{k-1}(z)] = \sum_{j=1}^{\mathcal{N}_0} \langle F_1 \widetilde{W} \Phi_j, \widetilde{W} \Phi_j \rangle + O(|z|^\epsilon). \quad (4.65)$$

The lemma comes from the following result (see [12], [30]) :

$$\langle F_1 \widetilde{W} \Phi_i, \widetilde{W} \Phi_j \rangle = \delta_{ij}. \quad (4.66)$$

\square

Now, we prove :

Lemma 4.11.

$$T_{33}(z) \approx 0.$$

Proof. Since $S_{23}(z) \approx 0$, it is easy to see that

$$T_{33}(z) \approx (-1)^{k+1} \sum_{\{\vec{\nu}\} \leq 1}^{(1)} \sum_{j, \{\vec{\nu}_1\} + j \leq 0}^{(-)} \frac{d}{dz} z_{\vec{\nu}} z_{\vec{\nu}_1} z^j \text{Tr} [U_1 R_{\vec{\nu}, 0} V T_{e; \vec{\nu}_1; j} U_2 S_3^{k-1}(z)]$$

We note that $T_{e;\vec{\nu};j} = \Pi_0 A_{e;\vec{\nu};j}$ with $A_{e;\vec{\nu};j}$ be a bounded operator in $\mathcal{L}(-1, s; -1, s)$, $s > 3$. So the result comes from Lemma 4.9. \square

We have the following estimate for $T_{32}(z)$:

Lemma 4.12.

$$T_{34}(z) \approx \frac{\delta_{\varsigma_{\kappa_0}, 1}}{z \ln z} \text{Tr} [R_1 V (\Pi_0 \widetilde{W} Q_e F_1 \widetilde{W} \Pi_{r, \kappa_0} + \Pi_{r, \kappa_0} \widetilde{W} Q_r F_1 \widetilde{W} \Pi_0) V].$$

Proof. We have

$$T_{34}(z) = (-1)^{k+1} \text{Tr} [(U_1 R_1 U_2 + \sum_{\{\vec{\nu}\} \leq 1}^{(1)} \frac{d}{dz} z_{\vec{\nu}} U_1 R_{\vec{\nu}, 0} U_2) (U_1 T_{er}(z) U_2) S_3^{k-1}(z)]. \quad (4.67)$$

First, we assume that $\varsigma_{\kappa_0} < 1$. In this case, $U_1 T_{er}(z) U_2 \approx 0$ and it follows that

$$T_{34}(z) \approx J_1(z) + J_2(z) + J_3(z), \quad (4.68)$$

where

$$\begin{aligned} J_1(z) &= (-1)^{k+1} \sum_{j=1}^{\kappa_0} \sum_{\{\vec{\nu}\} \leq 1}^{(1)} \frac{d}{dz} z_{\vec{\nu}} z_{\varsigma_j}^{-1} \text{Tr} [U_1 R_{\vec{\nu}, 0} V \Pi_0 \widetilde{W} Q_e F_1 \widetilde{W} \Pi_{r, j} U_2 S_3^{k-1}(z)]. \\ J_2(z) &= (-1)^{k+1} \sum_{j=1}^{\kappa_0} \sum_{\{\vec{\nu}\} \leq 1}^{(1)} \frac{d}{dz} z_{\vec{\nu}} z_{\varsigma_j}^{-1} \text{Tr} [U_1 R_{\vec{\nu}, 0} V \Pi_{r, j} \widetilde{W} Q_r F_1 \widetilde{W} \Pi_0 U_2 S_3^{k-1}(z)]. \\ J_3(z) &= (-1)^{k+1} \sum_{j=1}^{\kappa_0} \sum_{\{\vec{\nu}\} \leq 1}^{(1)} \sum_{\alpha, \beta, \vec{\nu}_1, l}^{+, 1} \frac{d}{dz} z_{\vec{\nu}} z_{\varsigma_j}^{-1} z_{\vec{\nu}_1} z^{|\beta|} (z_{\varsigma_j})^{-\alpha-\beta} z^l \\ &\quad \text{Tr} [U_1 R_{\vec{\nu}, 0} V T_{er; \vec{\nu}_1, \alpha, \beta, l, j} U_2 S_3^{k-1}(z)]. \end{aligned}$$

By Lemma 4.9, $J_1(z) = 0$. Now, let us study $J_2(z)$. As previously, the leading contribution is obtained when $\vec{\nu} = \varsigma_j$. So we have :

$$\begin{aligned} J_2(z) &\approx \frac{(-1)^{k+1}}{z} \sum_{j=1}^{\kappa_0} \varsigma_j \text{Tr} [U_1 R_{\varsigma_j, 0} V \Pi_{r, j} \widetilde{W} Q_r F_1 \widetilde{W} \Pi_0 U_2 S_3^{k-1}(z)] \\ &\approx \frac{(-1)^{k+1}}{z} \sum_{j=1}^{\kappa_0} \varsigma_j \text{Tr} [V \Pi_{r, j} \widetilde{W} Q_r F_1 \widetilde{W} \Pi_0 (V R_0)^{k-1} V R_{\varsigma_j, 0}], \end{aligned}$$

where we have used the cyclicity of the trace. Moreover, it is easy to see that Lemma 4.6 implies

$$\Pi_0 (V R_0)^{k-1} = (-1)^{k-1} \Pi_0, \quad \Pi_{r, j} (V R_0)^{k-1} = (-1)^{k-1} \Pi_{r, j}. \quad (4.69)$$

Using the first assertion of (4.69) and Lemma 4.9, we deduce

$$\begin{aligned} J_2(z) &\approx \frac{1}{z} \sum_{j=1}^{\kappa_0} \varsigma_j \text{Tr} [V \Pi_{r, j} \widetilde{W} Q_r F_1 \widetilde{W} \Pi_0 V R_{\varsigma_j, 0}] \\ &\approx 0. \end{aligned}$$

Finally, using the same arguments as in Lemma 4.8, we can show that $J_3(z) \approx 0$.

Now, let assume that $\varsigma_{k_0} = 1$. We follow the same strategy as in the case $\varsigma_{k_0} < 1$. It is easy to see that, in this case,

$$\begin{aligned} T_{34}(z) &\approx \frac{(-1)^{k+1}}{z \ln z} \operatorname{Tr} [U_1 R_1 V (\Pi_0 \widetilde{W} Q_e F_1 \widetilde{W} \Pi_{r, \kappa_0} + \Pi_{r, \kappa_0} \widetilde{W} Q_r F_1 \widetilde{W} \Pi_0) U_2 S_3^{k-1}(z)] \\ &\approx \frac{(-1)^{k+1}}{z \ln z} \operatorname{Tr} [R_1 V (\Pi_0 \widetilde{W} Q_e F_1 \widetilde{W} \Pi_{r, \kappa_0} + \Pi_{r, \kappa_0} \widetilde{W} Q_r F_1 \widetilde{W} \Pi_0) (V R_0)^{k-1} V]. \end{aligned}$$

Then, the lemma comes from (4.69). \square

Finally, $T_{31}(z)$ has been computed in the case where 0 is not an eigenvalue of P .

End of the proof of Theorem 4.4.

As a conclusion, it follows from the above discussion that in all the cases, one has

$$T(z) \approx -\frac{1}{z} \left(\sum_{j=1}^{\kappa_0} \varsigma_j m_{\varsigma_j} + \mathcal{N}_0 \right) + \frac{C}{z \ln z}, \quad (4.70)$$

for some constant C and for all z near 0 with $\Im z > 0$. Using the relation $T(z) = \overline{T(\bar{z})}$ for $\Im z < 0$, we deduce that C is real and (4.70) holds for z near 0 and $\Im z \neq 0$. Recall that the generalized residue is defined as

$$J_0 = -\frac{1}{2\pi i} \lim_{\delta \rightarrow 0} \lim_{\epsilon \rightarrow 0} \int_{\gamma(\delta, \epsilon)} T(z) dz$$

if the limit exists, where $\gamma(\delta, \epsilon)$ is positively oriented. We obtain from (4.70) that

$$-\frac{1}{2\pi i} \int_{\gamma(\delta, \epsilon)} T(z) dz = \sum_{j=1}^{\kappa_0} \varsigma_j m_{\varsigma_j} + \mathcal{N}_0 + O\left(\frac{\epsilon}{\delta}\right) + O\left(\frac{1}{|\log \delta|}\right),$$

uniformly in $0 < \epsilon < \delta$. Taking first the limit $\epsilon \rightarrow 0$, then the limit $\delta \rightarrow 0$, one derives $J_0 = \sum_{j=1}^{\kappa_0} \varsigma_j m_{\varsigma_j} + \mathcal{N}_0$. This proves Theorem 4.4. \square

5. PROOF OF LEVINSON'S THEOREM

In this section, we shall use Theorem 2.1 to prove a Levinson's theorem for potentials with critical decay. First, we have to verify that the conditions (2.1), and (2.4) – (2.9) hold. In the last section, we have seen that (2.8) is satisfied if $\rho_0 > \max(6, n+2)$.

Under the assumption (1.4), the Hamiltonian P_0 is positive, and it is well known that P and P_0 have no embedded positive eigenvalues, and the spectrum of P and P_0 is purely absolutely continuous on $]0, +\infty[$.

Using Theorem 3.7 with $N = 0$, one has for $s > 2$, z small with $\Im z > 0$,

$$\| \langle x \rangle^{-s} R(z) \langle x \rangle^{-s} \| \leq \frac{C}{|z|} \quad (5.1)$$

We deduce that the negative eigenvalues of P can not accumulate at 0. Indeed, let φ be an eigenfunction of P associated with an eigenvalue $\lambda \in [-\delta, 0[$, δ small. It is well known that φ decays exponentially. So, using (5.1) with $z = \lambda + i\epsilon$, we obtain

$$\frac{1}{\epsilon} \| \langle x \rangle^{-s} \varphi \| \leq \frac{C}{|\lambda + i\epsilon|} \| \langle x \rangle^s \varphi \|, \quad (5.2)$$

which gives the contradiction when $\epsilon \rightarrow 0$.

At least, for $\rho_0 > n$ and $k > \frac{n}{2}$, it is well known that (2.1) holds, (see for example, [25], Theorems 1.1 -1.2).

Now, let us prove the following elementary lemma in order to verify (2.6) :

Lemma 5.1. *Assume $\rho_0 > n$. Then, for every $f \in C_0^\infty(\mathbb{R})$, there exists $C > 0$, such that :*

$$| \operatorname{Tr} [(R(z) - R_0(z))f(P)] | \leq \frac{C}{|z|^2},$$

uniformly in $z \in \mathbb{C}$ with $|z|$ large and $z \notin \sigma(P)$.

Proof. For $z \notin \sigma(P)$, using the resolvent identity, we have :

$$\begin{aligned} (R(z) - R_0(z))f(P) &= -R(z)VR_0(z)f(P) \\ &= -R(z)VR(z)f(P) - R(z)VR_0(z)VR(z)f(P). \end{aligned}$$

Let $f_1 \in C_0^\infty$ such that $f_1 f = f$. Using the cyclicity of the trace, we obtain :

$$\operatorname{Tr} [(R(z) - R_0(z))f(P)] = (1) + (2),$$

with

$$(1) = -\operatorname{Tr}[(R(z)f_1(P)) (f_1(P)V) (R(z)f(P))], \quad (5.3)$$

$$\begin{aligned} (2) &= -\operatorname{Tr}[(R(z)f_1(P)) (f_1(P)\langle x \rangle^{-\frac{\rho_0}{2}}) \\ &\quad (\langle x \rangle^{\frac{\rho_0}{2}} V R_0(z) V \langle x \rangle^{\frac{\rho_0}{2}}) (\langle x \rangle^{-\frac{\rho_0}{2}} f_1(P)) (R(z)f(P))]. \end{aligned} \quad (5.4)$$

For $|z|$ large, $\Im z \neq 0$, we have ([27]),

$$\| \langle x \rangle^{\frac{\rho_0}{2}} V R_0(z) V \langle x \rangle^{\frac{\rho_0}{2}} \| = O(|z|^{-1/2}).$$

Moreover, by the spectral theorem,

$$\| R(z)f(P) \| = \frac{1}{\operatorname{dist}(z, \operatorname{Supp} f)} = O\left(\frac{1}{|z|}\right).$$

The lemma follows now from the fact that $\langle x \rangle^{-\frac{\rho_0}{2}} f_1(P)$ is a Hilbert-Schmidt operator and $f_1(P)V$ is of trace class. \square

In the same way, using Theorem 3.6 with $N = 0$, one has for $z \notin \sigma(P_0)$, $|z|$ small, and $s > 1$, $\rho_0 > n + 2$,

$$\begin{aligned} \operatorname{Tr}[R_0(z)(f(P) - f(P_0))] &= \operatorname{Tr}[\langle x \rangle^{-s}(R_0 + O(|z|^\epsilon))\langle x \rangle^{-s} \\ &\quad \langle x \rangle^s(f(P) - f(P_0))\langle x \rangle^s] \\ &= O(1). \end{aligned} \quad (5.5)$$

Then, the assumption (2.7) is satisfied.

In order to obtain a Levinson theorem, we have also to prove that $\xi'(\lambda)$ is integrable on $[0, 1]$. We have the following result :

Theorem 5.2. *Under the conditions of Theorem 4.4, for some $\epsilon_0 > 0$,*

$$\xi'(\lambda) = O(\lambda^{-1+\epsilon_0}), \quad \lambda \downarrow 0. \quad (5.6)$$

Proof. We shall prove that

$$\lambda^{1-\epsilon_0} \xi'(\lambda) \in L^\infty([0, 1]) = (L^1([0, 1]))', \quad (5.7)$$

or equivalently,

$$\exists C > 0, \forall \varphi \in C_0^\infty([0, 1]), \left| \int_{\mathbb{R}} \lambda^{1-\epsilon_0} \xi'(\lambda) \varphi(\lambda) d\lambda \right| \leq C \|\varphi\|_1. \quad (5.8)$$

Let us consider $g \in C_0^\infty([0, 1]; \mathbb{R})$ and $f \in C_0^\infty(\mathbb{R}; \mathbb{R})$ such that $f \equiv 1$ on $[0, 2]$, so we have $fg = g$. For $\lambda \in]0, 1[$ and $\epsilon > 0$, we have :

$$\begin{aligned} & \text{Tr} [R(\lambda + i\epsilon)f(P) - R_0(\lambda + i\epsilon)f(P_0)] \\ &= \text{Tr} [(R(\lambda + i\epsilon) - R_0(\lambda + i\epsilon))f(P)] + \text{Tr} [R_0(\lambda + i\epsilon) (f(P) - f(P_0))]. \end{aligned} \quad (5.9)$$

By the definition of the SSF,

$$\text{Tr} [R(\lambda + i\epsilon)f(P) - R_0(\lambda + i\epsilon)f(P_0)] = - \int_{\mathbb{R}} \xi(s) \frac{\partial f}{\partial s}(s, \lambda + i\epsilon) ds, \quad (5.10)$$

where we have set

$$f(s, \lambda + i\epsilon) = \frac{1}{s - \lambda - i\epsilon} f(s). \quad (5.11)$$

Thus, the left hand side of (5.9) is given by

$$(LHS) = \int_{\mathbb{R}} \xi(s) \frac{1}{(s - \lambda - i\epsilon)^2} f(s) ds - \int_{\mathbb{R}} \xi(s) \frac{1}{(s - \lambda - i\epsilon)} f'(s) ds. \quad (5.12)$$

Using (4.70) and (5.5), the right hand side of (5.9) satisfies for a suitable real constant C ,

$$(RHS) = -\frac{J_0}{\lambda + i\epsilon} + \frac{C}{(\lambda + i\epsilon) \log(\lambda + i\epsilon)} + r(\lambda, \epsilon), \quad (5.13)$$

where

$$r(\lambda, \epsilon) = O(|\lambda + i\epsilon|^{-1+\epsilon_0}) + \text{Tr} [R_0(\lambda + i\epsilon) (f(P) - f(P_0))] \quad (5.14)$$

$$= O(\lambda^{-1+\epsilon_0}), \quad \lambda \downarrow 0. \quad (5.15)$$

Thus, we have obtained that

$$r(\lambda, \epsilon) = \frac{J_0}{\lambda + i\epsilon} - \frac{C}{(\lambda + i\epsilon) \log(\lambda + i\epsilon)} + a(\lambda, \epsilon) - b(\lambda, \epsilon), \quad (5.16)$$

where

$$a(\lambda, \epsilon) = \int_{\mathbb{R}} \xi(s) \frac{1}{(s - \lambda - i\epsilon)^2} f(s) ds, \quad (5.17)$$

$$b(\lambda, \epsilon) = \int_{\mathbb{R}} \xi(s) \frac{1}{s - \lambda - i\epsilon} f'(s) ds. \quad (5.18)$$

We shall multiply (5.16) by $g(\lambda)$ and we shall integrate over λ . First, we remark that

$$\begin{aligned} \int_{\mathbb{R}} a(\lambda, \epsilon) g(\lambda) d\lambda &= - \int_{\mathbb{R}} \xi(s) f(s) \left(\int_{\mathbb{R}} \frac{\partial}{\partial \lambda} \left(\frac{1}{\lambda - s + i\epsilon} \right) g(\lambda) d\lambda \right) ds \\ &= \int_{\mathbb{R}} \xi(s) f(s) \left(\int_{\mathbb{R}} \frac{1}{\lambda - s + i\epsilon} g'(\lambda) d\lambda \right) ds. \end{aligned} \quad (5.19)$$

Thus,

$$\begin{aligned} \int_{\mathbb{R}} r(\lambda, \epsilon) g(\lambda) d\lambda &= \int_{\mathbb{R}} \frac{J_0}{\lambda + i\epsilon} g(\lambda) d\lambda - \int_{\mathbb{R}} \frac{C}{(\lambda + i\epsilon) \log(\lambda + i\epsilon)} g(\lambda) d\lambda \\ &+ \int_{\mathbb{R}} \xi(s) f(s) \left(\int_{\mathbb{R}} \frac{1}{\lambda - s + i\epsilon} g'(\lambda) d\lambda \right) ds \\ &+ \int_{\mathbb{R}} \xi(s) f'(s) \left(\int_{\mathbb{R}} \frac{1}{\lambda - s + i\epsilon} g(\lambda) d\lambda \right) ds. \end{aligned} \quad (5.20)$$

Taking the imaginary part in (5.20) and using that

$$\lim_{\epsilon \downarrow 0} \operatorname{Im} \frac{1}{\lambda - s + i\epsilon} = \pi \delta_{\lambda=s}, \quad (5.21)$$

we obtain easily

$$\lim_{\epsilon \downarrow 0} \operatorname{Im} \int_{\mathbb{R}} r(\lambda, \epsilon) g(\lambda) d\lambda = \pi \int_{\mathbb{R}} \xi(s) (fg)'(s) ds. \quad (5.22)$$

Now, we remark that :

$$\int_{\mathbb{R}} \xi(s) (fg)'(s) ds = \int_{\mathbb{R}} \xi(s) g'(s) ds = - \int_{\mathbb{R}} \xi'(s) g(s) ds, \quad (5.23)$$

since $\xi \in C^\infty([0, +\infty[)$. Then, we choose $g(s) = s^{1-\epsilon_0} \varphi(s)$ with $\varphi \in C_0^\infty([0, 1])$, and using (5.14), we have proved (5.8). \square

Now, we are able to prove a Levinson's theorem for critical potentials. To do it, we recall that if v satisfies (1.5), one has the high energy asymptotics, ([25], Theorem 1.2) :

$$\xi'(\lambda) \sim \sum_{j \geq 1} c_j \lambda^{\frac{n}{2}-j-1}, \quad \lambda \rightarrow +\infty. \quad (5.24)$$

This implies in particular that the condition (2.9) is satisfied. We also need the following asymptotics which comes from the functional calculus on pseudodifferential operators (see [25], Theorem 1.1): Let $\chi \in \mathcal{S}(\mathbb{R})$, with $\chi \equiv 1$ in a neighborhood of 0. We have :

$$\operatorname{Tr} \left[\chi\left(\frac{P}{R}\right) - \chi\left(\frac{P_0}{R}\right) \right] \sim \sum_{j \geq 1} \beta_j R^{\frac{n}{2}-j}, \quad R \rightarrow +\infty, \quad (5.25)$$

where the coefficients β_j are distributions on χ and is calculable, in principle, in terms of the symbols of P and P_0 .

The main result of this section is the following :

Theorem 5.3. *Assume $\rho_0 > \max(6, n+2)$. Then, we have :*

$$\int_0^\infty (\xi'(\lambda) - \sum_{j=1}^{\lfloor \frac{n}{2} \rfloor} c_j \lambda^{\lfloor \frac{n}{2} \rfloor - 1 - j}) d\lambda = -(\mathcal{N}_- + J_0) + \beta_{n/2}, \quad (5.26)$$

where $\beta_{n/2}$ depends only on n , v and V . If n is odd, $\beta_{n/2} = 0$. If n is even, $c_{n/2} = 0$.

Proof. Let $\chi \in C_0^\infty(\mathbb{R})$ such that $\chi(\lambda) \equiv 1$ in a neighborhood of 0. We use Theorem 2.1 with $f(\lambda) = \chi(\frac{\lambda}{R})$. The proof is divided in two steps.

Case 1. The dimension n is odd.

We first compute the second term on the left hand side of (2.10) as follows :

$$\begin{aligned} \int_0^\infty \chi\left(\frac{\lambda}{R}\right) \xi'(\lambda) d\lambda &= \int_0^\infty \chi\left(\frac{\lambda}{R}\right) \left[\xi'(\lambda) - \sum_{j=1}^{\lfloor \frac{n}{2} \rfloor} c_j \lambda^{\frac{n}{2}-1-j}\right] d\lambda + \sum_{j=1}^{\lfloor \frac{n}{2} \rfloor} c_j \int_0^\infty \chi\left(\frac{\lambda}{R}\right) \lambda^{\frac{n}{2}-1-j} d\lambda \\ &= \int_0^\infty \chi\left(\frac{\lambda}{R}\right) \left[\xi'(\lambda) - \sum_{j=1}^{\lfloor \frac{n}{2} \rfloor} c_j \lambda^{\frac{n}{2}-1-j}\right] d\lambda + \sum_{j=1}^{\lfloor \frac{n}{2} \rfloor} d_j R^{\frac{n}{2}-j}, \end{aligned}$$

with $d_j = c_j \int_0^\infty \chi(t) t^{\frac{n}{2}-1-j} dt$. Using (2.10) and (5.25), one has

$$\int_0^\infty \chi\left(\frac{\lambda}{R}\right) \left[\xi'(\lambda) - \sum_{j=1}^{\lfloor \frac{n}{2} \rfloor} c_j \lambda^{\frac{n}{2}-1-j}\right] d\lambda = -(\mathcal{N}_- + J_0) + \sum_{j=1}^{\lfloor \frac{n}{2} \rfloor} (\beta_j - d_j) R^{\frac{n}{2}-j} + O(R^{-\epsilon}) \quad (5.27)$$

where $\epsilon = 1 - \frac{n}{2} + \lfloor \frac{n}{2} \rfloor > 0$.

Now, let us study the left hand side of (5.27). We have :

$$\begin{aligned} &\int_0^\infty \chi\left(\frac{\lambda}{R}\right) \left[\xi'(\lambda) - \sum_{j=1}^{\lfloor \frac{n}{2} \rfloor} c_j \lambda^{\frac{n}{2}-1-j}\right] d\lambda \\ &= \int_0^1 \chi\left(\frac{\lambda}{R}\right) \left[\xi'(\lambda) - \sum_{j=1}^{\lfloor \frac{n}{2} \rfloor} c_j \lambda^{\frac{n}{2}-1-j}\right] d\lambda + \int_1^\infty \chi\left(\frac{\lambda}{R}\right) \left[\xi'(\lambda) - \sum_{j=1}^{\lfloor \frac{n}{2} \rfloor} c_j \lambda^{\frac{n}{2}-1-j}\right] d\lambda. \end{aligned}$$

By Theorem 5.2, $\xi'(\lambda)$ is integrable on $]0, 1[$, thus

$$\lim_{R \rightarrow +\infty} \int_0^1 \chi\left(\frac{\lambda}{R}\right) \left[\xi'(\lambda) - \sum_{j=1}^{\lfloor \frac{n}{2} \rfloor} c_j \lambda^{\frac{n}{2}-1-j}\right] d\lambda = \int_0^1 \left[\xi'(\lambda) - \sum_{j=1}^{\lfloor \frac{n}{2} \rfloor} c_j \lambda^{\frac{n}{2}-1-j}\right] d\lambda. \quad (5.28)$$

In the same way, (5.24) implies

$$\left| \xi'(\lambda) - \sum_{j=1}^{\lfloor \frac{n}{2} \rfloor} c_j \lambda^{\frac{n}{2}-1-j} \right| \leq C \lambda^{-\nu}, \quad (5.29)$$

with $\nu = 2 - \frac{n}{2} + \lfloor \frac{n}{2} \rfloor > 1$. Therefore,

$$\lim_{R \rightarrow +\infty} \int_1^\infty \chi\left(\frac{\lambda}{R}\right) \left[\xi'(\lambda) - \sum_{j=1}^{\lfloor \frac{n}{2} \rfloor} c_j \lambda^{\frac{n}{2}-1-j}\right] d\lambda = \int_1^\infty \left[\xi'(\lambda) - \sum_{j=1}^{\lfloor \frac{n}{2} \rfloor} c_j \lambda^{\frac{n}{2}-1-j}\right] d\lambda.$$

It follows that

$$\lim_{R \rightarrow +\infty} \int_0^\infty \chi\left(\frac{\lambda}{R}\right) \left[\xi'(\lambda) - \sum_{j=1}^{\lfloor \frac{n}{2} \rfloor} c_j \lambda^{\frac{n}{2}-1-j}\right] d\lambda = \int_0^\infty \left[\xi'(\lambda) - \sum_{j=1}^{\lfloor \frac{n}{2} \rfloor} c_j \lambda^{\frac{n}{2}-1-j}\right] d\lambda. \quad (5.30)$$

Taking $R \rightarrow +\infty$ in both sides of (5.27), we deduce that $\beta_j = d_j$. Thus, we get the Levinson's Theorem for odd dimension :

$$\int_0^\infty \left(\xi'(\lambda) - \sum_{j=1}^{\lfloor \frac{n}{2} \rfloor} c_j \lambda^{\lfloor \frac{n}{2} \rfloor - 1 - j}\right) d\lambda = -(\mathcal{N}_- + J_0). \quad (5.31)$$

Case 2: The dimension n is even.

First, let us study the case $n = 2$. Using Theorem 2.1, we have as previously :

$$\begin{aligned} & \int_0^1 \chi\left(\frac{\lambda}{R}\right) \xi'(\lambda) d\lambda + \int_1^\infty \chi\left(\frac{\lambda}{R}\right) \left[\xi'(\lambda) - \frac{c_1}{\lambda}\right] d\lambda \\ &= -(\mathcal{N}_- + J_0) + \beta_1 + O\left(\frac{1}{R}\right) - c_1 \int_1^{+\infty} \chi\left(\frac{\lambda}{R}\right) \frac{d\lambda}{\lambda}. \end{aligned} \quad (5.32)$$

We take $R \rightarrow +\infty$ in (5.32), and remarking that

$$\int_1^\infty \chi\left(\frac{\lambda}{R}\right) \frac{d\lambda}{\lambda} \sim \log R, \quad (5.33)$$

we deduce $c_1 = 0$ and we have obtained the Levinson theorem in dimension $n = 2$:

$$\int_0^{+\infty} \xi'(\lambda) d\lambda = -(\mathcal{N}_- + J_0) + \beta_1. \quad (5.34)$$

Now, assume $n \geq 4$. We write $n = 2p$ with $p \geq 2$. We define for $1 \leq j \leq p-1$,

$$d_j = c_j \int_0^\infty \chi(t) t^{p-1-j} dt. \quad (5.35)$$

As for the case $n = 2$, we obtain easily :

$$\begin{aligned} & \int_0^1 \chi\left(\frac{\lambda}{R}\right) \left[\xi'(\lambda) - \sum_{j=1}^{p-1} c_j \lambda^{p-1-j}\right] d\lambda + \int_1^\infty \chi\left(\frac{\lambda}{R}\right) \left[\xi'(\lambda) - \sum_{j=1}^p c_j \lambda^{p-1-j}\right] d\lambda \\ &= -(\mathcal{N}_- + J_0) + \sum_{j=1}^{p-1} (\beta_j - d_j) R^{p-j} + \beta_p + O\left(\frac{1}{R}\right) - c_p \int_1^\infty \chi\left(\frac{\lambda}{R}\right) \frac{d\lambda}{\lambda}. \end{aligned} \quad (5.36)$$

As previously, we take $R \rightarrow +\infty$ in (5.36) and we deduce that $\beta_j = d_j$ and $c_p = 0$. Therefore, we get the Levinson's theorem for even dimension:

$$\int_0^\infty \left(\xi'(\lambda) - \sum_{j=1}^{p-1} c_j \lambda^{p-1-j}\right) d\lambda = -(\mathcal{N}_- + J_0) + \beta_p. \quad (5.37)$$

□

Remark 5.4. Of course, the values of β_p appearing in the Levinson theorem are independent of the cutoff function χ . We can use the functional calculus on pseudodifferential operators ([9]) with the Hamiltonians $P_0 = -\Delta + v(x) - V(x)$ and $P = -\Delta + v(x)$, to compute β_p .

If $f \in C_0^\infty(\mathbb{R})$, we recall that $f(tP) - f(tP_0)$ is a \sqrt{t} -admissible operator, i.e

$$f(tP) - f(tP_0) = Op_{\sqrt{t}}(a_f(t)), \quad (5.38)$$

with

$$a_f(t) \sim \sum_{j \geq 1} a_{f,j} t^{\frac{j}{2}}. \quad (5.39)$$

The symbols $a_{f,j}$ are defined as

$$a_{f,j}(x, \xi) = \sum_{k=1}^{2j-1} (-1)^k d_{j,k}(x, \xi) f^{(k)}(\xi^2), \quad (5.40)$$

where the $d_{j,k}$ are universal polynomial functions on ξ which do not depend on f . For example, one has

$$d_{2,1}(x, \xi) = V(x), \quad d_{2,2}(x, \xi) = 0, \quad d_{2,3}(x, \xi) = 0. \quad (5.41)$$

$$d_{4,1}(x, \xi) = 0, \quad d_{4,2}(x, \xi) = 2v(x)V(x) - V^2(x), \quad (5.42)$$

$$d_{4,3}(x, \xi) = \sum_i \partial_i^2 V(x) \xi_i^2 - 2 \sum_{i < j} \partial_{ij} V(x) \xi_i \xi_j, \quad (5.43)$$

$$d_{4,k}(x, \xi) = 0 \quad \text{if } 4 \leq k \leq 7. \quad (5.44)$$

When $t \rightarrow 0$,

$$\text{Tr} [f(tP) - f(tP_0)] \sim t^{-\frac{n}{2}} \sum_{j \geq 1} \beta_j t^j, \quad (5.45)$$

with

$$\beta_j = (2\pi)^{-n} \sum_{k=1}^{4j-1} \int_{\mathbb{R}^n} \int_{\mathbb{R}^n} d_{2j,k}(x, \xi) f^{(k)}(\xi^2) dx d\xi. \quad (5.46)$$

As a consequence, if we set,

$$\gamma_n = \text{Vol} (S^{n-1}) = \frac{2\pi^{n/2}}{\Gamma(n/2)}, \quad (5.47)$$

where Γ is the well-known Gamma function, we can show :

- In dimension $n = 2$:

$$\beta_1 = \frac{1}{(2\pi)^2} \frac{\gamma_2}{2} \int_{\mathbb{R}^2} V(x) dx. \quad (5.48)$$

- In dimension $n = 4$:

$$\beta_2 = \frac{1}{(2\pi)^4} \frac{\gamma_4}{2} \int_{\mathbb{R}^4} (2v(x)V(x) - V^2(x)) dx. \quad (5.49)$$

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